

A LOCALLY DIVERGENCE-FREE INTERIOR PENALTY METHOD FOR TWO-DIMENSIONAL CURL-CURL PROBLEMS

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Abstract. An interior penalty method for certain two-dimensional curl-curl problems is investigated in this paper. This method computes the divergence-free part of the solution using locally divergence-free discontinuous P_1 vector fields on graded meshes. It has optimal order convergence (up to an arbitrarily small ϵ) for the source problem and the eigenproblem. Results of numerical experiments that corroborate the theoretical results are also presented.

Key words. curl-curl problem, Maxwell eigenproblem, locally divergence-free, interior penalty methods, graded meshes

AMS subject classification. 65N30, 65N15, 35Q60

1. Introduction. Let $\Omega \subset \mathbb{R}^2$ be a bounded polygonal domain. Consider the following weak curl-curl problem:

Find $\mathbf{u} \in H_0(\text{curl}; \Omega)$ such that

$$(1.1) \quad (\nabla \times \mathbf{u}, \nabla \times \mathbf{v}) + \alpha(\mathbf{u}, \mathbf{v}) = (\mathbf{f}, \mathbf{v}) \quad \forall \mathbf{v} \in H_0(\text{curl}; \Omega),$$

where α is a constant, (\cdot, \cdot) denotes the inner product of $L_2(\Omega)$ (or $[L_2(\Omega)]^2$),

$$H(\text{curl}; \Omega) = \left\{ \mathbf{v} = \begin{bmatrix} v_1 \\ v_2 \end{bmatrix} \in [L_2(\Omega)]^2 : \nabla \times \mathbf{v} = \frac{\partial v_2}{\partial x_1} - \frac{\partial v_1}{\partial x_2} \in L_2(\Omega) \right\}$$

and

$$H_0(\text{curl}; \Omega) = \left\{ \mathbf{v} = \begin{bmatrix} v_1 \\ v_2 \end{bmatrix} \in H(\text{curl}; \Omega) : \mathbf{n} \times \mathbf{v} = n_1 v_2 - n_2 v_1 = 0 \text{ on } \partial\Omega \right\},$$

with $\mathbf{n} = \begin{bmatrix} n_1 \\ n_2 \end{bmatrix}$ being the unit outer normal on $\partial\Omega$.

The curl-curl problem (1.1) is related to electromagnetic problems. For $\alpha \leq 0$, it is the weak form of the time-harmonic (frequency-domain) Maxwell equations. For $\alpha > 0$, it is related to the spatial problems appearing in implicit semi-discretizations of the time-dependent (time-domain) Maxwell equations.

It is well-known [24, 25, 5] that H^1 conforming nodal finite element methods can lead to a wrong solution of (1.1) if Ω is not convex. Many alternative approaches have been developed which include $H(\text{curl}; \Omega)$ conforming edge element methods [41, 42, 13, 33, 38, 40], H^1 conforming nodal finite methods with weighted regularization [27, 29], the singular complement/field method [32, 5], and interior penalty methods [43, 36, 37, 35, 34].

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In this paper we take a different approach. By the Helmholtz decomposition, the solution $\mathbf{u} \in H_0(\text{curl}; \Omega)$ of (1.1) can be written uniquely as

$$(1.2) \quad \mathbf{u} = \dot{\mathbf{u}} + \nabla \phi,$$

where $\dot{\mathbf{u}} \in H_0(\text{curl}; \Omega) \cap H(\text{div}^0; \Omega)$, $\phi \in H_0^1(\Omega)$,

$$H(\text{div}^0; \Omega) = \left\{ \mathbf{v} = \begin{bmatrix} v_1 \\ v_2 \end{bmatrix} \in [L_2(\Omega)]^2 : \nabla \cdot \mathbf{v} = \frac{\partial v_1}{\partial x_1} + \frac{\partial v_2}{\partial x_2} = 0 \right\}.$$

It is easy to see that ϕ satisfies

$$(1.3) \quad \alpha(\nabla \phi, \nabla \psi) = (\mathbf{f}, \nabla \psi) \quad \forall \psi \in H_0^1(\Omega).$$

Since the Poisson problem (1.3) (when $\alpha \neq 0$) can be solved by many standard methods under the assumption that $\mathbf{f} \in H(\text{div}; \Omega)$, we will focus on the divergence-free part $\dot{\mathbf{u}}$, which satisfies

$$(1.4) \quad (\nabla \times \dot{\mathbf{u}}, \nabla \times \mathbf{v}) + \alpha(\dot{\mathbf{u}}, \mathbf{v}) = (\mathbf{f}, \mathbf{v}) \quad \forall \mathbf{v} \in H_0(\text{curl}; \Omega) \cap H(\text{div}^0; \Omega).$$

We shall refer to (1.4) as the weak form of the reduced curl-curl problem and assume in the case where $\alpha \leq 0$ that (1.4) has a unique solution (i.e. $-\alpha$ is not a Maxwell eigenvalue). Note that the strong form of the reduced curl-curl problem is given by

$$(1.5) \quad \nabla \times (\nabla \times \dot{\mathbf{u}}) + \alpha \dot{\mathbf{u}} = Q\mathbf{f},$$

where $Q : [L_2(\Omega)]^2 \rightarrow H(\text{div}^0; \Omega)$ is the orthogonal projection.

Remark 1.1. The strong form (1.5) and the well-posedness of (1.4) imply immediately

$$(1.6) \quad \|\nabla \times \dot{\mathbf{u}}\|_{L_2(\Omega)} \leq C_{\Omega, \alpha} \|Q\mathbf{f}\|_{L_2(\Omega)} \leq C_{\Omega, \alpha} \|\mathbf{f}\|_{L_2(\Omega)}.$$

Here and below we use C (with or without subscripts) to denote a generic positive constant that can take different values at different occurrences.

The advantage of working with (1.4) is that it behaves like an elliptic problem, unlike the full curl-curl problem (1.1). In particular, the solution $\dot{\mathbf{u}}$ of (1.4) enjoys elliptic regularity under the assumption that $\mathbf{f} \in L_2(\Omega)$, which greatly simplifies the duality argument needed for the analysis.

In [15] we developed a numerical scheme (in the case where $\alpha \leq 0$) for (1.4) using the Crouzeix-Raviart nonconforming P_1 vector fields [30]. In this paper we investigate an interior penalty method that avoids the difficulties involved in the construction of local bases for the Crouzeix-Raviart vector fields at the cost of a larger number of unknowns. These two methods are very closely related, which simplifies the analysis of the new scheme.

Since the solution operator of the interior penalty method converges uniformly to the solution operator of (1.4) with respect to the norm $\|\cdot\|_{L_2(\Omega)}$, classical spectral approximation theory [7] implies that we can also apply the interior penalty method to the following Maxwell eigenproblem:

Find $(\lambda, \dot{\mathbf{u}}) \in \mathbb{R} \times H_0(\text{curl}; \Omega) \cap H(\text{div}^0; \Omega)$ such that

$$(1.7) \quad (\nabla \times \dot{\mathbf{u}}, \nabla \times \mathbf{v}) = \lambda(\dot{\mathbf{u}}, \mathbf{v}) \quad \forall \mathbf{v} \in H_0(\text{curl}; \Omega) \cap H(\text{div}^0; \Omega).$$

In particular, the spectral approximation based on our interior penalty method is free of spurious eigenvalues.

The rest of the paper is organized as follows. We introduce the finite element space and graded meshes in Section 2. The interior penalty method and some preliminary estimates are presented in Section 3. The convergence analysis for the source problem (1.1) is carried out in Section 4. Application to the eigenproblem (1.7) is discussed in Section 5, followed by numerical results in Section 6, and we end with some concluding remarks in Section 7.

2. Finite Element Spaces and Graded Meshes. In this section we will introduce discontinuous finite element spaces defined on meshes graded around the corners c_1, \dots, c_L of Ω .

We assume that the triangulation \mathcal{T}_h of Ω satisfies the following condition:

$$(2.1) \quad h_T \approx h\Phi_\mu(T) \quad \forall T \in \mathcal{T}_h,$$

where $h_T = \text{diam } T$, h is a mesh parameter proportional to $\max_{T \in \mathcal{T}_h} h_T$, and $\mu = (\mu_1, \dots, \mu_L)$ is the vector containing the grading parameters. The constants in the equivalence (2.1) are independent of h and depend in general on the minimum angle of the mesh. The weight $\Phi_\mu(T)$ is defined by

$$(2.2) \quad \Phi_\mu(T) = \prod_{\ell=1}^L |c_\ell - c_T|^{1-\mu_\ell},$$

where c_T is the center of T . Note that

$$(2.3) \quad h_T \leq Ch.$$

The choices of the grading parameters are dictated by the singularities [6, 28, 26] of the solution $\hat{\mathbf{u}}$ of (1.4). In order to recover optimal convergence rates in both the energy norm and the L_2 norm (cf. [15]), we take

$$(2.4) \quad \begin{aligned} \mu_\ell &= 1 & \text{if } \omega_\ell \leq \frac{\pi}{2}, \\ \mu_\ell &< \frac{\pi}{2\omega_\ell} & \text{if } \omega_\ell > \frac{\pi}{2}, \end{aligned}$$

where ω_ℓ is the interior angle at the corner c_ℓ . The construction of \mathcal{T}_h satisfying (2.1) can be found for example in [31, 4, 1, 14]. Note that \mathcal{T}_h can be constructed so that it satisfies a minimum angle condition for any fixed choice of the grading parameters.

Remark 2.1. The choice of the grading parameters in (2.4) indicates that grading is needed at any corner where the angle is larger than $\pi/2$. This is different from problems involving the Laplace operator where grading is only needed at re-entrant corners and is due to the fact that the singularities of the curl-curl operator are one order worse than the singularities of the Laplace operator.

The finite element space V_h is defined by

$$V_h = \{\mathbf{v} \in [L_2(\Omega)]^2 : \mathbf{v}_T = \mathbf{v}|_T \in [P_1(T)]^2 \quad \text{and} \quad \nabla \cdot \mathbf{v}_T = 0 \quad \forall T \in \mathcal{T}_h\}.$$

For any $s > \frac{1}{2}$ there is a natural weak interpolation operator $\Pi_T : [H^s(T)]^2 \rightarrow [P_1(T)]^2$ defined by

$$(2.5) \quad (\Pi_T \boldsymbol{\zeta})(m_{e_i}) = \frac{1}{|e_i|} \int_{e_i} \boldsymbol{\zeta} \, ds \quad \text{for } i = 1, 2, 3,$$

where m_{e_i} is the midpoint of the edge e_i of T and $|e|$ denotes the length of an edge e .

It follows immediately from (2.5), the midpoint rule and Green's theorem that

$$(2.6) \quad \int_T \nabla \times (\Pi_T \zeta) dx = \int_T \nabla \times \zeta dx \quad \forall \zeta \in [H^s(T)]^2,$$

$$(2.7) \quad \int_T \nabla \cdot (\Pi_T \zeta) dx = \int_T \nabla \cdot \zeta dx \quad \forall \zeta \in [H^s(T)]^2.$$

Furthermore, given $s \in (1/2, 2]$, we have the following interpolation error estimates [30]:

$$(2.8) \quad \|\zeta - \Pi_T \zeta\|_{L_2(T)} + h_T^{\min(s,1)} |\zeta - \Pi_T \zeta|_{H^{\min(s,1)}(T)} \leq C_T h_T^s |\zeta|_{H^s(T)}$$

for all $\zeta \in [H^s(T)]^2$, where the positive constant C_T depends on the minimum angle of T (and also on s when s is close to $1/2$).

Since the solution $\hat{\mathbf{u}}$ of the reduced curl-curl problem belongs to $[H^s(\Omega)]^2$ for some $s > \frac{1}{2}$ (cf. [40]), we can define a global interpolant of $\hat{\mathbf{u}}$ by

$$(\Pi_h \hat{\mathbf{u}})_T = \Pi_T \hat{\mathbf{u}}_T \quad \forall T \in \mathcal{T}_h,$$

where $\hat{\mathbf{u}}_T = \hat{\mathbf{u}}|_T$. It follows from (2.7) and $\nabla \cdot \hat{\mathbf{u}} = 0$ that $\Pi_h \hat{\mathbf{u}} \in V_h$.

Remark 2.2. Note that $\Pi_h \hat{\mathbf{u}}$ is the same interpolation operator used in [15] and $\Pi_h \hat{\mathbf{u}}$ belongs to the space of locally divergence-free Crouzeix-Raviart nonconforming P_1 vector fields [30].

Since the vector fields in V_h are discontinuous, their jumps across the edges of \mathcal{T}_h play an important role in the development of interior penalty methods. We will denote by \mathcal{E}_h (resp. \mathcal{E}_h^i) the set of the edges (resp. interior edges) of \mathcal{T}_h . Let $e \in \mathcal{E}_h^i$ be shared by the two triangles $T_1, T_2 \in \mathcal{T}_h$ (cf. Figure 2.1) and \mathbf{n}_1 (resp. \mathbf{n}_2) be the unit normal of e pointing towards the outside of T_1 (resp. T_2). We define, on e ,

$$(2.9a) \quad \llbracket \mathbf{n} \times \mathbf{v} \rrbracket = \mathbf{n}_1 \times \mathbf{v}_{T_1}|_e + \mathbf{n}_2 \times \mathbf{v}_{T_2}|_e,$$

$$(2.9b) \quad \llbracket \mathbf{n} \cdot \mathbf{v} \rrbracket = \mathbf{n}_1 \cdot \mathbf{v}_{T_1}|_e + \mathbf{n}_2 \cdot \mathbf{v}_{T_2}|_e.$$

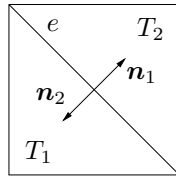


FIG. 2.1. Triangles and normals in the definitions of $\llbracket \mathbf{n} \times \mathbf{v} \rrbracket$ and $\llbracket \mathbf{n} \cdot \mathbf{v} \rrbracket$

For an edge e along $\partial\Omega$, we take \mathbf{n}_e to be the unit normal of e pointing towards the outside of Ω and define

$$(2.10) \quad \llbracket \mathbf{n} \times \mathbf{v} \rrbracket = \mathbf{n}_e \times \mathbf{v}|_e.$$

We will also denote the piecewise defined curl and div operators by $\nabla_h \times$ and $\nabla_h \cdot$, i.e.,

$$(\nabla_h \times \mathbf{v})|_T = \nabla \times \mathbf{v}_T \quad \text{and} \quad (\nabla_h \cdot \mathbf{v})|_T = \nabla \cdot \mathbf{v}_T.$$

3. Discretization and Preliminary Error Estimates. In this section we present the interior penalty method for the reduced curl-curl problem and some preliminary estimates.

The discrete problem is to find $\hat{\mathbf{u}}_h \in V_h$ such that

$$(3.1) \quad a_h(\hat{\mathbf{u}}_h, \mathbf{v}) = (\mathbf{f}, \mathbf{v}) \quad \forall \mathbf{v} \in V_h,$$

where

$$(3.2) \quad \begin{aligned} a_h(\mathbf{w}, \mathbf{v}) &= (\nabla_h \times \mathbf{w}, \nabla_h \times \mathbf{v}) + \alpha(\mathbf{w}, \mathbf{v}) \\ &+ \sum_{e \in \mathcal{E}_h} \frac{[\Phi_\mu(e)]^2}{|e|} \int_e \llbracket \mathbf{n} \times \mathbf{w} \rrbracket \llbracket \mathbf{n} \times \mathbf{v} \rrbracket ds \\ &+ \sum_{e \in \mathcal{E}_h^i} \frac{[\Phi_\mu(e)]^2}{|e|} \int_e \llbracket \mathbf{n} \cdot \mathbf{w} \rrbracket \llbracket \mathbf{n} \cdot \mathbf{v} \rrbracket ds \\ &+ h^{-2} \sum_{e \in \mathcal{E}_h} \frac{1}{|e|} \int_e (\Pi_e^0 \llbracket \mathbf{n} \times \mathbf{w} \rrbracket) (\Pi_e^0 \llbracket \mathbf{n} \times \mathbf{v} \rrbracket) ds \\ &+ h^{-2} \sum_{e \in \mathcal{E}_h^i} \frac{1}{|e|} \int_e (\Pi_e^0 \llbracket \mathbf{n} \cdot \mathbf{w} \rrbracket) (\Pi_e^0 \llbracket \mathbf{n} \cdot \mathbf{v} \rrbracket) ds, \end{aligned}$$

where Π_e^0 is the orthogonal projection from $L_2(e)$ onto $P_0(e)$ (the space of constant functions on e). The weight $\Phi_\mu(e)$ in (3.2) is defined by

$$\Phi_\mu(e) = \Pi_{\ell=1}^L |c_\ell - m_e|^{1-\mu_\ell},$$

where m_e is the midpoint of e .

Remark 3.1. The weight $\Phi_\mu(e)$ is closely related to the weight $\Phi_\mu(T)$ in (2.2). In fact, we have

$$(3.3) \quad \Phi_\mu(e) \approx \Phi_\mu(T) \quad \text{if } e \subset \partial T,$$

where the constants in the equivalence are independent of h and depend in general on the minimum angle of the mesh. The inclusion of $\Phi_\mu(e)$ in (3.2) is crucial for the derivation of optimal convergence rates on graded meshes.

Remark 3.2. The weak problem (1.4) for the reduced curl-curl problem can be written as

$$a_h(\hat{\mathbf{u}}, \mathbf{v}) = (\mathbf{f}, \mathbf{v}) \quad \forall \mathbf{v} \in H_0(\text{curl}; \Omega) \cap H(\text{div}^0; \Omega).$$

Remark 3.3. For nonconforming Crouzeix-Raviart P_1 vector fields \mathbf{w} and \mathbf{v} , we have

$$\begin{aligned} a_h(\mathbf{w}, \mathbf{v}) &= (\nabla_h \times \mathbf{w}, \nabla_h \times \mathbf{v}) + \alpha(\mathbf{w}, \mathbf{v}) \\ &+ \sum_{e \in \mathcal{E}_h} \frac{[\Phi_\mu(e)]^2}{|e|} \int_e \llbracket \mathbf{n} \times \mathbf{w} \rrbracket \llbracket \mathbf{n} \times \mathbf{v} \rrbracket ds \\ &+ \sum_{e \in \mathcal{E}_h^i} \frac{[\Phi_\mu(e)]^2}{|e|} \int_e \llbracket \mathbf{n} \cdot \mathbf{w} \rrbracket \llbracket \mathbf{n} \cdot \mathbf{v} \rrbracket ds, \end{aligned}$$

which is the variational form used in [15]. The two additional sums in (3.2) compensate for the lack of weak continuity for the vector fields in V_h .

Remark 3.4. The weak over-penalization causes the discrete system to become more ill-conditioned. For example, the condition number grows at the rate of $O(h^{-4})$ on quasi-uniform meshes. However, there exists a simple block-diagonal preconditioner that can reduce the growth of the condition number to $O(h^{-2})$ on quasi-uniform meshes and $O(h^{-2}(1 + |\ln h|))$ on graded meshes [16, 17].

Let the mesh-dependent energy norm $\|\cdot\|_h$ be defined by

$$(3.4) \quad \begin{aligned} \|\mathbf{v}\|_h^2 &= \|\nabla_h \times \mathbf{v}\|_{L_2(\Omega)}^2 + \|\mathbf{v}\|_{L_2(\Omega)}^2 \\ &+ \sum_{e \in \mathcal{E}_h} \frac{[\Phi_\mu(e)]^2}{|e|} \|[\mathbf{n} \times \mathbf{v}]\|_{L_2(e)}^2 + \sum_{e \in \mathcal{E}_h^i} \frac{[\Phi_\mu(e)]^2}{|e|} \|[\mathbf{n} \cdot \mathbf{v}]\|_{L_2(e)}^2 \\ &+ h^{-2} \left(\sum_{e \in \mathcal{E}_h} \frac{1}{|e|} \|\Pi_e^0[\mathbf{n} \times \mathbf{v}]\|_{L_2(e)}^2 + \sum_{e \in \mathcal{E}_h^i} \frac{1}{|e|} \|\Pi_e^0[\mathbf{n} \cdot \mathbf{v}]\|_{L_2(e)}^2 \right). \end{aligned}$$

Note that

$$(3.5) \quad \|\mathbf{v}\|_{L_2(\Omega)} \leq \|\mathbf{v}\|_h \quad \forall \mathbf{v} \in H_0(\text{curl}; \Omega) \cap H(\text{div}^0; \Omega) + V_h,$$

and from (3.2) and (3.4),

$$(3.6) \quad |a_h(\mathbf{w}, \mathbf{v})| \leq (|\alpha| + 1) \|\mathbf{w}\|_h \|\mathbf{v}\|_h \quad \forall \mathbf{w}, \mathbf{v} \in [H_0(\text{curl}; \Omega) \cap H(\text{div}^0; \Omega)] + V_h.$$

For $\alpha > 0$, $a_h(\cdot, \cdot)$ is coercive with respect to $\|\cdot\|_h$, i.e.,

$$(3.7) \quad a_h(\mathbf{v}, \mathbf{v}) \geq \min(1, \alpha) \|\mathbf{v}\|_h^2 \quad \forall \mathbf{v} \in H_0(\text{curl}; \Omega) \cap H(\text{div}^0; \Omega) + V_h.$$

In this case the discrete problem is well-posed and we have the following abstract error estimate.

LEMMA 3.5. *Let α be positive, $\beta = \min(1, \alpha)$, $\hat{\mathbf{u}}$ be the solution of (1.4) and $\hat{\mathbf{u}}_h$ satisfy (3.1). It holds that*

$$(3.8) \quad \|\hat{\mathbf{u}} - \hat{\mathbf{u}}_h\|_h \leq \left(\frac{1 + \alpha + \beta}{\beta} \right) \inf_{\mathbf{v} \in V_h} \|\hat{\mathbf{u}} - \mathbf{v}\|_h + \frac{1}{\beta} \max_{\mathbf{w} \in V_h \setminus \{\mathbf{0}\}} \frac{a_h(\hat{\mathbf{u}} - \hat{\mathbf{u}}_h, \mathbf{w})}{\|\mathbf{w}\|_h}.$$

Proof. Let $\mathbf{v} \in V_h$ be arbitrary. It follows from (3.6), (3.7) and the triangle inequality that

$$\begin{aligned} \|\hat{\mathbf{u}} - \hat{\mathbf{u}}_h\|_h &\leq \|\hat{\mathbf{u}} - \mathbf{v}\|_h + \|\mathbf{v} - \hat{\mathbf{u}}_h\|_h \\ &\leq \|\hat{\mathbf{u}} - \mathbf{v}\|_h + \frac{1}{\beta} \max_{\mathbf{w} \in V_h \setminus \{\mathbf{0}\}} \frac{a_h(\mathbf{v} - \hat{\mathbf{u}}_h, \mathbf{w})}{\|\mathbf{w}\|_h} \\ &\leq \left(\frac{1 + \alpha + \beta}{\beta} \right) \|\hat{\mathbf{u}} - \mathbf{v}\|_h + \frac{1}{\beta} \max_{\mathbf{w} \in V_h \setminus \{\mathbf{0}\}} \frac{a_h(\hat{\mathbf{u}} - \hat{\mathbf{u}}_h, \mathbf{w})}{\|\mathbf{w}\|_h} \end{aligned}$$

which implies (3.8). \square

For $\alpha \leq 0$, the following Gårding (in)equality holds:

$$(3.9) \quad a_h(\mathbf{v}, \mathbf{v}) + (|\alpha| + 1)(\mathbf{v}, \mathbf{v}) = \|\mathbf{v}\|_h^2 \quad \forall \mathbf{v} \in [H_0(\text{curl}; \Omega) \cap H(\text{div}^0; \Omega)] + V_h.$$

In this case the discrete problem is indefinite and the following lemma provides an abstract error estimate for the scheme (3.1) under the assumption that it has a solution.

LEMMA 3.6. *Let $\hat{\mathbf{u}} \in H_0(\text{curl}; \Omega) \cap H(\text{div}^0; \Omega)$ satisfy (1.4) and $\hat{\mathbf{u}}_h$ be a solution of (3.1). It holds that*

$$(3.10) \quad \|\hat{\mathbf{u}} - \hat{\mathbf{u}}_h\|_h \leq (2|\alpha| + 3) \inf_{\mathbf{v} \in V_h} \|\hat{\mathbf{u}} - \mathbf{v}\|_h + \max_{\mathbf{w} \in V_h \setminus \{\mathbf{0}\}} \frac{a_h(\hat{\mathbf{u}} - \hat{\mathbf{u}}_h, \mathbf{w})}{\|\mathbf{w}\|_h} + (|\alpha| + 1) \|\hat{\mathbf{u}} - \hat{\mathbf{u}}_h\|_{L_2(\Omega)}.$$

Proof. It follows from (3.5) and (3.9) that, for $\mathbf{v} \in V_h \setminus \{\mathbf{0}\}$,

$$(3.11) \quad \begin{aligned} \|\mathbf{v}\|_h &\leq \frac{a_h(\mathbf{v}, \mathbf{v})}{\|\mathbf{v}\|_h} + (|\alpha| + 1) \frac{(\mathbf{v}, \mathbf{v})}{\|\mathbf{v}\|_h} \\ &\leq \max_{\mathbf{w} \in V_h \setminus \{\mathbf{0}\}} \frac{a_h(\mathbf{v}, \mathbf{w})}{\|\mathbf{w}\|_h} + (|\alpha| + 1) \|\mathbf{v}\|_{L_2(\Omega)}. \end{aligned}$$

Let $\mathbf{v} \in V_h$ be arbitrary. We find, using (3.5), (3.6), (3.11) and the triangle inequality,

$$\begin{aligned} \|\hat{\mathbf{u}} - \hat{\mathbf{u}}_h\|_h &\leq \|\hat{\mathbf{u}} - \mathbf{v}\|_h + \|\mathbf{v} - \hat{\mathbf{u}}_h\|_h \\ &\leq \|\hat{\mathbf{u}} - \mathbf{v}\|_h + \max_{\mathbf{w} \in V_h \setminus \{\mathbf{0}\}} \frac{a_h(\mathbf{v} - \hat{\mathbf{u}}_h, \mathbf{w})}{\|\mathbf{w}\|_h} + (|\alpha| + 1) \|\mathbf{v} - \hat{\mathbf{u}}_h\|_{L_2(\Omega)} \\ &\leq (2|\alpha| + 3) \|\hat{\mathbf{u}} - \mathbf{v}\|_h + \max_{\mathbf{w} \in V_h \setminus \{\mathbf{0}\}} \frac{a_h(\hat{\mathbf{u}} - \hat{\mathbf{u}}_h, \mathbf{w})}{\|\mathbf{w}\|_h} + (|\alpha| + 1) \|\hat{\mathbf{u}} - \hat{\mathbf{u}}_h\|_{L_2(\Omega)}, \end{aligned}$$

which implies (3.10). \square

From here on we consider α to be fixed and drop the dependence on α in our estimates.

Remark 3.7. *The first term on the right-hand side of (3.8) and (3.10) measures the approximation property of V_h with respect to the norm $\|\cdot\|_h$. The second term measures the consistency error. The third term on the right-hand side of (3.10) addresses the indefiniteness of the reduced curl-curl problem when $\alpha \leq 0$.*

As mentioned in Remark 2.2, the interpolation operator Π_h is also the one used in [15]. Therefore in our analysis we can use the following two results from that paper (cf. Lemma 5.1 and Lemma 5.2 of [15]), which were obtained using (2.4), (2.8) and a representation of $\hat{\mathbf{u}}$ as the sum of a regular part and a singular part.

LEMMA 3.8. *Let $\hat{\mathbf{u}} \in H_0(\text{curl}; \Omega) \cap H(\text{div}^0; \Omega)$ be the solution of (1.4). We have the following interpolation error estimate:*

$$(3.12) \quad \|\hat{\mathbf{u}} - \Pi_h \hat{\mathbf{u}}\|_{L_2(\Omega)} \leq C_\epsilon h^{2-\epsilon} \|\mathbf{f}\|_{L_2(\Omega)} \quad \text{for any } \epsilon > 0.$$

LEMMA 3.9. *Let $\hat{\mathbf{u}} \in H_0(\text{curl}; \Omega) \cap H(\text{div}^0; \Omega)$ be the solution of (1.4). We have the following interpolation error estimate:*

$$(3.13) \quad \sum_{e \in \mathcal{E}_h} \frac{[\Phi_\mu(e)]^2}{|e|} \|\hat{\mathbf{u}} - \Pi_h \hat{\mathbf{u}}\|_{L_2(e)}^2 \leq C_\epsilon h^{2-\epsilon} \|\mathbf{f}\|_{L_2(\Omega)}^2$$

for any $\epsilon > 0$, where $[[\dot{\mathbf{u}} - \Pi_h \dot{\mathbf{u}}]]$ is the jump of $\dot{\mathbf{u}} - \Pi_h \dot{\mathbf{u}}$ across the interior edges of \mathcal{T}_h and $[[\dot{\mathbf{u}} - \Pi_h \dot{\mathbf{u}}]] = \dot{\mathbf{u}} - \Pi_h \dot{\mathbf{u}}$ on the boundary edges of \mathcal{T}_h .

LEMMA 3.10. *It holds that*

$$\sum_{e \in \mathcal{E}_h} |e| [\Phi_\mu(e)]^{-2} \|\eta - \bar{\eta}_{T_e}\|_{L_2(e)}^2 \leq Ch^2 |\eta|_{H^1(\Omega)}^2 \quad \forall \eta \in H^1(\Omega),$$

where

$$(3.14) \quad \bar{\eta}_{T_e} = \frac{1}{|T_e|} \int_{T_e} \eta \, dx$$

is the mean of η over T_e , one of the triangles in \mathcal{T}_h that has e as an edge.

Proof. This is the consequence of (2.1), (3.3), the trace theorem (with scaling) and a standard interpolation error estimate [22, 18]:

$$\begin{aligned} & \sum_{e \in \mathcal{E}_h} |e| [\Phi_\mu(e)]^{-2} \|\eta - \bar{\eta}_{T_e}\|_{L_2(e)}^2 \\ & \leq C \sum_{e \in \mathcal{E}_h} [\Phi_\mu(T)]^{-2} (\|\eta - \bar{\eta}_{T_e}\|_{L_2(T_e)}^2 + h_T^2 |\eta - \bar{\eta}_{T_e}|_{H^1(T_e)}^2) \\ & \leq C \sum_{e \in \mathcal{E}_h} [\Phi_\mu(T)]^{-2} h_T^2 |\eta|_{H^1(T_e)}^2 \leq Ch^2 |\eta|_{H^1(\Omega)}^2. \end{aligned}$$

□

Recall that Q is the L_2 orthogonal projection operator onto $H(\operatorname{div}^0; \Omega)$. The following result will be useful in addressing the consistency error caused by the appearance of Q in (1.5).

LEMMA 3.11. *The following estimate holds:*

$$(3.15) \quad \|\mathbf{v} - Q\mathbf{v}\|_{L_2(\Omega)} \leq Ch \|\mathbf{v}\|_h \quad \forall \mathbf{v} \in [H_0(\operatorname{curl}; \Omega) \cap H(\operatorname{div}^0; \Omega)] + V_h.$$

Proof. Let $\mathbf{v} \in [H_0(\operatorname{curl}; \Omega) \cap H(\operatorname{div}^0; \Omega)] + V_h$ be arbitrary. Since $\mathbf{v} - Q\mathbf{v}$ belongs to $\nabla H_0^1(\Omega)$, the orthogonal complement of $H(\operatorname{div}^0; \Omega)$ in $[L_2(\Omega)]^2$, we have, by duality,

$$(3.16) \quad \|\mathbf{v} - Q\mathbf{v}\|_{L_2(\Omega)} = \sup_{\eta \in H_0^1(\Omega) \setminus \{0\}} \frac{(\mathbf{v} - Q\mathbf{v}, \nabla \eta)}{\|\nabla \eta\|_{L_2(\Omega)}} = \sup_{\eta \in H_0^1(\Omega) \setminus \{0\}} \frac{(\mathbf{v}, \nabla \eta)}{\|\nabla \eta\|_{L_2(\Omega)}}.$$

Let $\eta \in H_0^1(\Omega)$ be arbitrary. Since $\nabla \cdot \mathbf{v} = 0$ on each triangle $T \in \mathcal{T}_h$, we find using integration by parts and the fact that $\eta|_{\partial\Omega} = 0$,

$$(3.17) \quad (\mathbf{v}, \nabla \eta) = \sum_{e \in \mathcal{E}_h^i} \int_e (\eta - \bar{\eta}_{T_e}) [[\mathbf{n} \cdot \mathbf{v}]] \, ds + \sum_{e \in \mathcal{E}_h^i} \int_e \bar{\eta}_{T_e} (\Pi_e^0 [[\mathbf{n} \cdot \mathbf{v}]]) \, ds = S_1 + S_2,$$

where $\bar{\eta}_{T_e}$ is defined in (3.14).

By the Cauchy-Schwarz inequality, (3.4) and Lemma 3.10, we have

$$(3.18) \quad \begin{aligned} S_1 & \leq \left[\sum_{e \in \mathcal{E}_h^i} |e| [\Phi_\mu(e)]^{-2} \|\eta - \bar{\eta}_{T_e}\|_{L_2(e)}^2 \right]^{1/2} \left[\sum_{e \in \mathcal{E}_h^i} \frac{[\Phi_\mu(e)]^2}{|e|} \|[[\mathbf{n} \cdot \mathbf{v}]]\|_{L_2(e)}^2 \right]^{1/2} \\ & \leq Ch \|\nabla \eta\|_{L_2(\Omega)} \|\mathbf{v}\|_h. \end{aligned}$$

On the other hand, from (3.4), the Cauchy-Schwarz inequality, and a Poincaré-Friedrichs inequality, we find

$$\begin{aligned}
(3.19) \quad S_2 &\leq \sum_{e \in \mathcal{E}_h^i} (|e|^{1/2} \|\bar{\eta}_{T_e}\|_{L_2(e)}) (|e|^{-1/2} \|\Pi_e^0[\mathbf{n} \cdot \mathbf{v}]\|_{L_2(e)}) \\
&\leq Ch \left[\sum_{e \in \mathcal{E}_h^i} \|\bar{\eta}_{T_e}\|_{L_2(T_e)}^2 \right]^{1/2} \left[h^{-2} \sum_{e \in \mathcal{E}_h^i} \frac{1}{|e|} \|\Pi_e^0[\mathbf{n} \cdot \mathbf{v}]\|_{L_2(e)}^2 \right]^{1/2} \\
&\leq Ch \|\eta\|_{L_2(\Omega)} \|\mathbf{v}\|_h \leq Ch \|\nabla \eta\|_{L_2(\Omega)} \|\mathbf{v}\|_h.
\end{aligned}$$

Here we have also used the simple fact that, if e is an edge of a triangle T , then

$$(3.20) \quad |e| \|q\|_{L_2(e)}^2 \leq C_T \|q\|_{L_2(T)}^2 \quad \text{for any constant function } q,$$

where the positive constant C_T depends only on the shape regularity of T .

The estimate (3.15) follows from (3.16)–(3.19). \square

4. Convergence Analysis. We begin with three lemmas that provide estimates to the three terms on the right-hand side of (3.8) and (3.10).

LEMMA 4.1. *Let $\hat{\mathbf{u}} \in H_0(\text{curl}; \Omega) \cap H(\text{div}^0; \Omega)$ be the solution of (1.4). It holds that*

$$(4.1) \quad \inf_{\mathbf{v} \in V_h} \|\hat{\mathbf{u}} - \mathbf{v}\|_h \leq \|\hat{\mathbf{u}} - \Pi_h \hat{\mathbf{u}}\|_h \leq C_\epsilon h^{1-\epsilon} \|\mathbf{f}\|_{L_2(\Omega)}$$

for any $\epsilon > 0$.

Proof. Since $\Pi_e^0[\mathbf{n} \times (\hat{\mathbf{u}} - \Pi_h \hat{\mathbf{u}})] = 0$ for all $e \in \mathcal{E}_h$ and $\Pi_e^0[\mathbf{n} \cdot (\hat{\mathbf{u}} - \Pi_h \hat{\mathbf{u}})] = 0$ for all $e \in \mathcal{E}_h^i$, we have

$$\begin{aligned}
(4.2) \quad \|\hat{\mathbf{u}} - \Pi_h \hat{\mathbf{u}}\|_h^2 &= \|\nabla_h \times (\hat{\mathbf{u}} - \Pi_h \hat{\mathbf{u}})\|_{L_2(\Omega)}^2 + \|\hat{\mathbf{u}} - \Pi_h \hat{\mathbf{u}}\|_{L_2(\Omega)}^2 \\
&+ \sum_{e \in \mathcal{E}_h} \frac{[\Phi_\mu(e)]^2}{|e|} \|\llbracket \mathbf{n} \times (\hat{\mathbf{u}} - \Pi_h \hat{\mathbf{u}}) \rrbracket\|_{L_2(e)}^2 \\
&+ \sum_{e \in \mathcal{E}_h^i} \frac{[\Phi_\mu(e)]^2}{|e|} \|\llbracket \mathbf{n} \cdot (\hat{\mathbf{u}} - \Pi_h \hat{\mathbf{u}}) \rrbracket\|_{L_2(e)}^2.
\end{aligned}$$

The second term on the right-hand side of (4.2) has been estimated in Lemma 3.8, and the third and fourth terms can be estimated using Lemma 3.9. Therefore it only remains to estimate the first term.

Observe that (2.6) implies

$$(4.3) \quad \nabla_h \times (\Pi_h \hat{\mathbf{u}}) = \Pi_h^0(\nabla \times \hat{\mathbf{u}}),$$

where Π_h^0 is the orthogonal projection from $L_2(\Omega)$ onto the space of piecewise constant functions with respect to \mathcal{T}_h . It then follows from (1.6), (2.3), (4.3) and a standard interpolation error estimate [22, 18] that

$$\begin{aligned}
(4.4) \quad \|\nabla_h \times (\hat{\mathbf{u}} - \Pi_h \hat{\mathbf{u}})\|_{L_2(\Omega)}^2 &= \|\nabla \times \hat{\mathbf{u}} - \Pi_h^0(\nabla \times \hat{\mathbf{u}})\|_{L_2(\Omega)}^2 \\
&\leq Ch^2 \|\nabla \times \hat{\mathbf{u}}\|_{H^1(\Omega)}^2 \leq Ch^2 \|\mathbf{f}\|_{L_2(\Omega)}^2.
\end{aligned}$$

The estimate (4.1) follows from (4.2), (4.4) and Lemmas 3.8–3.9. \square

LEMMA 4.2. *Let $\hat{\mathbf{u}} \in H_0(\text{curl}; \Omega) \cap H(\text{div}^0; \Omega)$ be the solution of (1.4) and $\hat{\mathbf{u}}_h \in V_h$ satisfy (3.1). It holds that*

$$(4.5) \quad \max_{\mathbf{w} \in V_h \setminus \{\mathbf{0}\}} \frac{a_h(\hat{\mathbf{u}} - \hat{\mathbf{u}}_h, \mathbf{w})}{\|\mathbf{w}\|_h} \leq Ch \|\mathbf{f}\|_{L_2(\Omega)}.$$

Proof. Let $\mathbf{w} \in V_h$ be arbitrary. Using integration by parts, the discrete problem (3.1), the strong form of the reduced curl-curl problem (1.5), and the fact that $\nabla_h \cdot \mathbf{w} = 0$, we find

$$(4.6) \quad \begin{aligned} a_h(\hat{\mathbf{u}} - \hat{\mathbf{u}}_h, \mathbf{w}) &= a_h(\hat{\mathbf{u}}, \mathbf{w}) - (\mathbf{f}, \mathbf{w}) \\ &= (Q\mathbf{f} - \mathbf{f}, \mathbf{w}) + \sum_{T \in \mathcal{T}_h} \int_{\partial T} (\nabla \times \hat{\mathbf{u}})(\mathbf{n}_T \times \mathbf{w}_T) ds, \end{aligned}$$

where $w_T = w|_T$ and \mathbf{n}_T is the unit outer normal along ∂T . We see from (4.6) that there are two sources for the inconsistency of the scheme defined by (3.1), namely the projection Q that appears in (1.5) and the discontinuity of the vector fields in V_h .

The equation (4.6) can be rewritten as

$$(4.7) \quad \begin{aligned} a_h(\hat{\mathbf{u}} - \hat{\mathbf{u}}_h, \mathbf{w}) &= (\mathbf{f}, Q\mathbf{w} - \mathbf{w}) + \sum_{e \in \mathcal{E}_h} \int_e (\nabla \times \hat{\mathbf{u}})[\mathbf{n} \times \mathbf{w}] ds \\ &= (\mathbf{f}, Q\mathbf{w} - \mathbf{w}) + \sum_{e \in \mathcal{E}_h} \int_e (\nabla \times \hat{\mathbf{u}} - \overline{(\nabla \times \hat{\mathbf{u}})}_{T_e})[\mathbf{n} \times \mathbf{w}] ds \\ &\quad + \sum_{e \in \mathcal{E}_h} \int_e \overline{(\nabla \times \hat{\mathbf{u}})}_{T_e} (\Pi_e^0[\mathbf{n} \times \mathbf{w}]) ds \end{aligned}$$

where $\overline{(\nabla \times \hat{\mathbf{u}})}_{T_e}$ is the mean of $\nabla \times \hat{\mathbf{u}}$ on one of the triangles $T_e \in \mathcal{T}_h$ that has e as an edge. The first two terms on the right-hand side of (4.7) satisfy the estimate

$$(4.8) \quad (\mathbf{f}, Q\mathbf{w} - \mathbf{w}) + \sum_{e \in \mathcal{E}_h} \int_e (\nabla \times \hat{\mathbf{u}} - \overline{(\nabla \times \hat{\mathbf{u}})}_{T_e})[\mathbf{n} \times \mathbf{w}] ds \leq Ch \|\mathbf{f}\|_{L_2(\Omega)} \|\mathbf{w}\|_h.$$

The derivation of (4.8), which is based on Lemma 3.10, Lemma 3.11 and (1.6), can be found in the proof of Lemma 6.2 in [15].

Using the Cauchy-Schwarz inequality, (1.6), (3.4) and (3.20), the third term on the right-hand side of (4.7) can be estimated as follows:

$$(4.9) \quad \begin{aligned} &\sum_{e \in \mathcal{E}_h} \int_e \overline{(\nabla \times \hat{\mathbf{u}})}_{T_e} (\Pi_e^0[\mathbf{n} \times \mathbf{w}]) ds \\ &\leq \sum_{e \in \mathcal{E}_h} (|e|^{1/2} \|\overline{(\nabla \times \hat{\mathbf{u}})}_{T_e}\|_{L_2(e)}) (|e|^{-1/2} \|\Pi_e^0[\mathbf{n} \times \mathbf{w}]\|_{L_2(e)}) \\ &\leq Ch \left(\sum_{e \in \mathcal{E}_h} \|\overline{(\nabla \times \hat{\mathbf{u}})}_{T_e}\|_{L_2(T_e)}^2 \right)^{1/2} \left(h^{-2} \sum_{e \in \mathcal{E}_h} \frac{1}{|e|} \|\Pi_e^0[\mathbf{n} \times \mathbf{w}]\|_{L_2(e)}^2 \right)^{1/2} \\ &\leq Ch \|\nabla \times \hat{\mathbf{u}}\|_{L_2(\Omega)} \|\mathbf{w}\|_h \leq Ch \|\mathbf{f}\|_{L_2(\Omega)} \|\mathbf{w}\|_h. \end{aligned}$$

The estimate (4.5) follows from (4.7)–(4.9). \square

LEMMA 4.3. *Let $\hat{\mathbf{u}} \in H_0(\text{curl}; \Omega) \cap H(\text{div}^0; \Omega)$ be the solution of (1.4) and $\hat{\mathbf{u}}_h \in V_h$ satisfy (3.1). It holds that*

$$(4.10) \quad \|\hat{\mathbf{u}} - \hat{\mathbf{u}}_h\|_{L_2(\Omega)} \leq C_\epsilon (h^{2-\epsilon} \|\mathbf{f}\|_{L_2(\Omega)} + h^{1-\epsilon} \|\hat{\mathbf{u}} - \hat{\mathbf{u}}_h\|_h)$$

for any $\epsilon > 0$.

Proof. We use a duality argument. Let $\hat{\mathbf{z}} \in H_0(\text{curl}; \Omega) \cap H(\text{div}^0; \Omega)$ satisfy the reduced curl-curl problem

$$(4.11) \quad a_h(\mathbf{v}, \hat{\mathbf{z}}) = (\mathbf{v}, \hat{\mathbf{u}} - \hat{\mathbf{u}}_h) \quad \forall \mathbf{v} \in H_0(\text{curl}; \Omega) \cap H(\text{div}^0; \Omega).$$

The strong form of (4.11) is

$$(4.12) \quad \nabla \times (\nabla \times \hat{\mathbf{z}}) + \alpha \hat{\mathbf{z}} = Q(\hat{\mathbf{u}} - \hat{\mathbf{u}}_h),$$

and we have the following analog of (1.6):

$$(4.13) \quad \|\nabla \times \hat{\mathbf{z}}\|_{H^1(\Omega)} \leq C_\Omega \|\hat{\mathbf{u}} - \hat{\mathbf{u}}_h\|_{L_2(\Omega)}.$$

It follows from (4.12) and integration by parts that

$$(4.14) \quad a_h(\hat{\mathbf{u}}_h, \hat{\mathbf{z}}) = (\hat{\mathbf{u}}_h, Q(\hat{\mathbf{u}} - \hat{\mathbf{u}}_h)) + \sum_{e \in \mathcal{E}_h} \int_e (\nabla \times \hat{\mathbf{z}}) \llbracket \mathbf{n} \times \hat{\mathbf{u}}_h \rrbracket ds.$$

From (4.11) and (4.14), we have

$$(4.15) \quad \begin{aligned} \|\hat{\mathbf{u}} - \hat{\mathbf{u}}_h\|_{L_2(\Omega)}^2 &= (\hat{\mathbf{u}}, \hat{\mathbf{u}} - \hat{\mathbf{u}}_h) - (\hat{\mathbf{u}}_h, \hat{\mathbf{u}} - \hat{\mathbf{u}}_h) \\ &= a_h(\hat{\mathbf{u}} - \hat{\mathbf{u}}_h, \hat{\mathbf{z}}) + a_h(\hat{\mathbf{u}}_h, \hat{\mathbf{z}}) - (\hat{\mathbf{u}}_h, \hat{\mathbf{u}} - \hat{\mathbf{u}}_h) \\ &= a_h(\hat{\mathbf{u}} - \hat{\mathbf{u}}_h, \hat{\mathbf{z}}) - (\hat{\mathbf{u}}_h, (I - Q)(\hat{\mathbf{u}} - \hat{\mathbf{u}}_h)) \\ &\quad + \sum_{e \in \mathcal{E}_h} \int_e (\nabla \times \hat{\mathbf{z}}) \llbracket \mathbf{n} \times \hat{\mathbf{u}}_h \rrbracket ds. \end{aligned}$$

We will estimate the three terms on the right-hand side of (4.15) separately.

Using (4.7) and the fact that $\llbracket \mathbf{n} \times (\Pi_h \hat{\mathbf{z}}) \rrbracket$ vanishes at the midpoints of all $e \in \mathcal{E}_h$, we can rewrite the first term as

$$(4.16) \quad \begin{aligned} a_h(\hat{\mathbf{u}} - \hat{\mathbf{u}}_h, \hat{\mathbf{z}}) &= a_h(\hat{\mathbf{u}} - \hat{\mathbf{u}}_h, \hat{\mathbf{z}} - \Pi_h \hat{\mathbf{z}}) + a_h(\hat{\mathbf{u}} - \hat{\mathbf{u}}_h, \Pi_h \hat{\mathbf{z}}) \\ &= a_h(\hat{\mathbf{u}} - \hat{\mathbf{u}}_h, \hat{\mathbf{z}} - \Pi_h \hat{\mathbf{z}}) + (\mathbf{f}, Q \Pi_h \hat{\mathbf{z}} - \Pi_h \hat{\mathbf{z}}) \\ &\quad + \sum_{e \in \mathcal{E}_h} \int_e (\nabla \times \hat{\mathbf{u}} - \overline{(\nabla \times \hat{\mathbf{u}})}_{T_e}) \llbracket \mathbf{n} \times (\Pi_h \hat{\mathbf{z}}) \rrbracket ds, \end{aligned}$$

where $\overline{(\nabla \times \hat{\mathbf{u}})}_{T_e}$ is the mean of $\nabla \times \hat{\mathbf{u}}$ on one of the triangles $T_e \in \mathcal{T}_h$ that has e as an edge. Using the expression (4.16), the following estimate was obtained in the proof of Lemma 6.5 of [15]:

$$(4.17) \quad a_h(\hat{\mathbf{u}} - \hat{\mathbf{u}}_h, \hat{\mathbf{z}}) \leq C_\epsilon (h^{2-\epsilon} \|\mathbf{f}\|_{L_2(\Omega)} + h^{1-\epsilon} \|\hat{\mathbf{u}} - \hat{\mathbf{u}}_h\|_h) \|\hat{\mathbf{u}} - \hat{\mathbf{u}}_h\|_{L_2(\Omega)}.$$

Since $(I - Q)\hat{\mathbf{u}} = \mathbf{0}$, the second term on the right-hand side of (4.15) can be estimated by Lemma 3.11:

$$(4.18) \quad -(\hat{\mathbf{u}}_h, (I - Q)(\hat{\mathbf{u}} - \hat{\mathbf{u}}_h)) = (\hat{\mathbf{u}} - \hat{\mathbf{u}}_h, (I - Q)(\hat{\mathbf{u}} - \hat{\mathbf{u}}_h))$$

$$\leq C \|\dot{\mathbf{u}} - \dot{\mathbf{u}}_h\|_{L_2(\Omega)} (h \|\dot{\mathbf{u}} - \dot{\mathbf{u}}_h\|_h).$$

Finally we estimate the third term on the right-hand side of (4.15). We have

$$(4.19) \quad \sum_{e \in \mathcal{E}_h} \int_e (\nabla \times \dot{\mathbf{z}}) [\mathbf{n} \times \dot{\mathbf{u}}_h] ds = \sum_{e \in \mathcal{E}_h} \int_e (\nabla \times \dot{\mathbf{z}} - \overline{(\nabla \times \dot{\mathbf{z}})}_{T_e}) [\mathbf{n} \times \dot{\mathbf{u}}_h] ds \\ + \sum_{e \in \mathcal{E}_h} \int_e \overline{(\nabla \times \dot{\mathbf{z}})}_{T_e} (\Pi_e^0 [\mathbf{n} \times \dot{\mathbf{u}}_h]) ds,$$

where $\overline{(\nabla \times \dot{\mathbf{z}})}_{T_e}$ is the mean of $\nabla \times \dot{\mathbf{z}}$ on a triangle $T_e \in \mathcal{T}_h$ that has e as an edge. The following estimate can be found in the proof of Lemma 6.5 of [15]:

$$(4.20) \quad \sum_{e \in \mathcal{E}_h} \int_e (\nabla \times \dot{\mathbf{z}} - \overline{(\nabla \times \dot{\mathbf{z}})}_{T_e}) [\mathbf{n} \times \dot{\mathbf{u}}_h] ds \leq Ch \|\dot{\mathbf{u}} - \dot{\mathbf{u}}_h\|_{L_2(\Omega)} \|\dot{\mathbf{u}} - \dot{\mathbf{u}}_h\|_h.$$

On the other hand, we obtain by the Cauchy-Schwarz inequality, (3.4), (3.20) and (4.13),

$$(4.21) \quad \sum_{e \in \mathcal{E}_h} \int_e \overline{(\nabla \times \dot{\mathbf{z}})}_{T_e} (\Pi_e^0 [\mathbf{n} \times \dot{\mathbf{u}}_h]) ds \\ = \sum_{e \in \mathcal{E}_h} \int_e \overline{(\nabla \times \dot{\mathbf{z}})}_{T_e} (\Pi_e^0 [\mathbf{n} \times (\dot{\mathbf{u}}_h - \dot{\mathbf{u}})]) ds \\ \leq Ch \|\nabla \times \dot{\mathbf{z}}\|_{L_2(\Omega)} \|\dot{\mathbf{u}} - \dot{\mathbf{u}}_h\|_h \leq Ch \|\dot{\mathbf{u}} - \dot{\mathbf{u}}_h\|_{L_2(\Omega)} \|\dot{\mathbf{u}} - \dot{\mathbf{u}}_h\|_h.$$

The estimate (4.10) follows from (4.15) and (4.17)–(4.21). \square

In the case where $\alpha > 0$, the following theorem is an immediate consequence of Lemma 3.5 and Lemmas 4.1–4.3.

THEOREM 4.4. *Let α be positive. The following estimates hold for the solution $\dot{\mathbf{u}}_h$ of (3.1):*

$$(4.22) \quad \|\dot{\mathbf{u}} - \dot{\mathbf{u}}_h\|_h \leq C_\epsilon h^{1-\epsilon} \|\mathbf{f}\|_{L_2(\Omega)} \quad \text{for any } \epsilon > 0,$$

$$(4.23) \quad \|\dot{\mathbf{u}} - \dot{\mathbf{u}}_h\|_{L_2(\Omega)} \leq C_\epsilon h^{2-\epsilon} \|\mathbf{f}\|_{L_2(\Omega)} \quad \text{for any } \epsilon > 0.$$

In the case where $\alpha \leq 0$, we have the following convergence theorem for the scheme (3.1).

THEOREM 4.5. *Assume that $-\alpha \geq 0$ is not a Maxwell eigenvalue. There exists a positive number h_* such that the discrete problem (3.1) is uniquely solvable for all $h \leq h_*$, in which case the following discretization error estimates are valid:*

$$(4.24) \quad \|\dot{\mathbf{u}} - \dot{\mathbf{u}}_h\|_h \leq C_\epsilon h^{1-\epsilon} \|\mathbf{f}\|_{L_2(\Omega)} \quad \text{for any } \epsilon > 0,$$

$$(4.25) \quad \|\dot{\mathbf{u}} - \dot{\mathbf{u}}_h\|_{L_2(\Omega)} \leq C_\epsilon h^{2-\epsilon} \|\mathbf{f}\|_{L_2(\Omega)} \quad \text{for any } \epsilon > 0.$$

Proof. We follow the approach of Schatz [44] for indefinite problems. Assuming $\dot{\mathbf{u}}_h$ satisfies (3.1), it follows from (3.10) and Lemmas 4.1–4.3 that

$$(4.26) \quad \|\dot{\mathbf{u}} - \dot{\mathbf{u}}_h\|_h \leq C_\epsilon h^{1-\epsilon} (\|\mathbf{f}\|_{L_2(\Omega)} + \|\dot{\mathbf{u}} - \dot{\mathbf{u}}_h\|_h) \quad \text{for any } \epsilon > 0.$$

By choosing an $\epsilon_* > 0$, we deduce from (4.26) that, for $h \leq h_* = 1/(2C_{\epsilon_*})^{1/(1-\epsilon_*)}$,

$$\begin{aligned} \|\hat{\mathbf{u}} - \hat{\mathbf{u}}_h\|_h &\leq C_{\epsilon_*} h^{1-\epsilon_*} (\|\mathbf{f}\|_{L_2(\Omega)} + \|\hat{\mathbf{u}} - \hat{\mathbf{u}}_h\|_h) \\ &\leq C_{\epsilon_*} h^{1-\epsilon_*} \|\mathbf{f}\|_{L_2(\Omega)} + C_{\epsilon_*} h_*^{1-\epsilon_*} \|\hat{\mathbf{u}} - \hat{\mathbf{u}}_h\|_h \\ &\leq C_{\epsilon_*} h^{1-\epsilon_*} \|\mathbf{f}\|_{L_2(\Omega)} + \frac{1}{2} \|\hat{\mathbf{u}} - \hat{\mathbf{u}}_h\|_h, \end{aligned}$$

and hence

$$(4.27) \quad \|\hat{\mathbf{u}} - \hat{\mathbf{u}}_h\|_h \leq 2C_{\epsilon_*} h^{1-\epsilon_*} \|\mathbf{f}\|_{L_2(\Omega)}.$$

Therefore, any solution $\hat{\mathbf{z}}_h \in V_h$ of the homogeneous discrete problem

$$(4.28) \quad a_h(\hat{\mathbf{z}}_h, \mathbf{v}) = 0 \quad \forall \mathbf{v} \in V_h,$$

which corresponds to the special case where $\mathbf{f} = \mathbf{0} = \hat{\mathbf{z}}$, will satisfy the following special case of (4.27):

$$\|\hat{\mathbf{z}}_h\|_h \leq 0.$$

Hence the only solution of (4.28) is the trivial solution and the discrete problem (3.1) is uniquely solvable for $h \leq h_*$.

The energy error estimate (4.24) now follows (4.27), and the L_2 error estimate (4.25) follows from Lemma 4.3 and (4.24). \square

5. Application to the Maxwell Eigenproblem. Given any $\mathbf{f} \in [L_2(\Omega)]^2$, we define $T\mathbf{f} \in H_0(\text{curl}; \Omega) \cap H(\text{div}^0; \Omega)$ by the condition that

$$(5.1) \quad (\nabla \times (T\mathbf{f}), \nabla \times \mathbf{v}) + (T\mathbf{f}, \mathbf{v}) = (\mathbf{f}, \mathbf{v}) \quad \forall \mathbf{v} \in H_0(\text{curl}; \Omega) \cap H(\text{div}^0; \Omega).$$

Clearly T is a bounded linear operator from $[L_2(\Omega)]^2$ into $H_0(\text{curl}; \Omega) \cap H(\text{div}^0; \Omega)$. Since $H_0(\text{curl}; \Omega) \cap H(\text{div}^0; \Omega)$ is a compact subspace of $[L_2(\Omega)]^2$ (cf. [23]), the operator $T : [L_2(\Omega)]^2 \rightarrow [L_2(\Omega)]^2$ is symmetric, positive and compact. Moreover, $(\lambda, \hat{\mathbf{u}})$ satisfy the Maxwell eigenproblem (1.7) if and only if

$$(5.2) \quad T\hat{\mathbf{u}} = \frac{1}{1+\lambda}\hat{\mathbf{u}}.$$

Similarly, let $T_h\mathbf{f} \in V_h$ be defined by the condition

$$(5.3) \quad a_{h,1}(T_h\mathbf{f}, \mathbf{v}) = (\mathbf{f}, \mathbf{v}) \quad \forall \mathbf{v} \in V_h,$$

where $a_{h,1}(\cdot, \cdot)$ is the bilinear form $a_h(\cdot, \cdot)$ in (3.2) with $\alpha = 1$, then

$$(5.4) \quad T_h\hat{\mathbf{u}}_h = \frac{1}{1+\lambda_h}\hat{\mathbf{u}}_h$$

is equivalent to

$$(5.5) \quad a_{h,0}(\hat{\mathbf{u}}_h, \mathbf{v}) = \lambda_h(\hat{\mathbf{u}}_h, \mathbf{v}) \quad \forall \mathbf{v} \in V_h,$$

where $a_{h,0}(\cdot, \cdot)$ is the bilinear form $a_h(\cdot, \cdot)$ in (3.2) with $\alpha = 0$.

From Theorem 4.4 we have

$$(5.6) \quad \|(T - T_h)\mathbf{f}\|_{L_2(\Omega)} \leq C_\epsilon h^{2-\epsilon} \|\mathbf{f}\|_{L_2(\Omega)}$$

for all $\mathbf{f} \in [L_2(\Omega)]^2$ and

$$(5.7) \quad \|(T - T_h)\mathbf{f}\|_h \leq C_\epsilon h^{1-\epsilon} \|\mathbf{f}\|_{L_2(\Omega)} \leq C_\epsilon h^{1-\epsilon} \|\mathbf{f}\|_h$$

for all $\mathbf{f} \in H_0(\text{curl}; \Omega) \cap H(\text{div}^0; \Omega) + V_h$. Thus the symmetric finite rank operator T_h converges uniformly to the symmetric positive compact operator T as $h \downarrow 0$, which implies that the classical theory of spectral approximation [39, 7] can be applied to the approximation of the eigenvalues and eigenfunctions of T by the eigenvalues and eigenfunctions of T_h . Hence we have the following result through the connections between (1.7) and (5.2) and between (5.5) and (5.4).

THEOREM 5.1. *Let $0 \leq \lambda_1 \leq \lambda_2 \leq \dots$ be the eigenvalues of (1.7), $\lambda = \lambda_j = \lambda_{j+1} = \dots = \lambda_{j+m-1}$ be an eigenvalue with multiplicity m , and $V_\lambda \subset H_0(\text{curl}; \Omega) \cap H(\text{div}^0; \Omega)$ be the corresponding m dimensional eigenspace. Let $\lambda_{h,1} \leq \lambda_{h,2} \leq \dots$ be the eigenvalues of (5.5). Then, as $h \downarrow 0$, we have*

$$(5.8) \quad |\lambda_{h,\ell} - \lambda| \leq C_{\lambda,\epsilon} h^{2-\epsilon} \quad \ell = j, \dots, j+m-1.$$

Furthermore, if $V_{h,\lambda} \subset V_h$ is the space spanned by the eigenfunctions corresponding to $\lambda_{h,j}, \dots, \lambda_{h,j+m-1}$, then the gap between V_λ and $V_{h,\lambda}$ goes to zero at the rate of $C_\epsilon h^{2-\epsilon}$ in the L_2 norm and at the rate of $C_\epsilon h^{1-\epsilon}$ in the norm $\|\cdot\|_h$.

Proof. Let E_λ (resp. $E_{h,\lambda}$) be the L_2 orthogonal projection onto V_λ (resp. $V_{h,\lambda}$). It follows from (5.6)–(5.7) and the classical theory of spectral approximation of compact operators [39, 7] that

$$(5.9) \quad \|(E_\lambda - E_{h,\lambda})\mathbf{w}\|_{L_2(\Omega)} \leq C_\epsilon h^{2-\epsilon} \|\mathbf{w}\|_{L_2(\Omega)}$$

for all $\mathbf{w} \in [L_2(\Omega)]^2$ and

$$(5.10) \quad \|(E_\lambda - E_{h,\lambda})\mathbf{w}\|_h \leq C_\epsilon h^{1-\epsilon} \|\mathbf{w}\|_h$$

for all $\mathbf{w} \in H_0(\text{curl}; \Omega) \cap H(\text{div}^0; \Omega) + V_h$.

Let $\mathbf{x} \in V_{h,\lambda}$ be a unit eigenfunction of $\lambda_{h,\ell}$ for some ℓ between j and $j+m-1$. Then we have

$$(5.11) \quad (T_h \mathbf{x}, \mathbf{x}) = \mu_{h,\ell} \quad \text{and} \quad \|\mathbf{x}\|_{L_2(\Omega)} = 1,$$

where (cf. (5.4))

$$(5.12) \quad \mu_{h,\ell} = (1 + \lambda_{h,\ell})^{-1}.$$

Let $\hat{\mathbf{x}} = E_\lambda \mathbf{x}$ and $\hat{\mathbf{y}} = \hat{\mathbf{x}} - \mathbf{x}$. Then $\hat{\mathbf{x}}$ and $\hat{\mathbf{y}}$ are orthogonal with respect to both the L_2 inner product and the inner product $(T \cdot, \cdot)$, and the estimate (5.9) implies that

$$(5.13) \quad \|\hat{\mathbf{y}}\|_{L_2(\Omega)} = \|(E_\lambda - E_{h,\lambda})\mathbf{x}\|_{L_2(\Omega)} \leq C_\epsilon h^{2-\epsilon}.$$

In particular, it follows from Pythagoras' theorem with respect to the L_2 inner product that

$$(5.14) \quad 1 - \|\hat{\mathbf{x}}\|_{L_2(\Omega)}^2 = \|\mathbf{x}\|_{L_2(\Omega)}^2 - \|\hat{\mathbf{x}}\|_{L_2(\Omega)}^2 = \|\hat{\mathbf{y}}\|_{L_2(\Omega)}^2 \leq C_\epsilon h^{4-\epsilon}.$$

Let $\hat{\mathbf{e}} = \hat{\mathbf{x}} / \|\hat{\mathbf{x}}\|_{L_2(\Omega)}$. Then $\hat{\mathbf{e}}$ is a unit eigenfunction of T corresponding to the eigenvalue (cf. (5.2))

$$(5.15) \quad \mu = (1 + \lambda)^{-1},$$

and we have

$$(5.16) \quad |\mu - \mu_{h,\ell}| = |(T\hat{\mathbf{e}}, \hat{\mathbf{e}}) - (T_h\mathbf{x}, \mathbf{x})| \leq |((T - T_h)\mathbf{x}, \mathbf{x})| + |(T\hat{\mathbf{e}}, \hat{\mathbf{e}}) - (T\mathbf{x}, \mathbf{x})|.$$

From (5.6) we have

$$(5.17) \quad |((T - T_h)\mathbf{x}, \mathbf{x})| \leq C_\epsilon h^{2-\epsilon},$$

and it follows from Pythagoras' theorem with respect to the inner product $(T\cdot, \cdot)$ that

$$(5.18) \quad \begin{aligned} |(T\hat{\mathbf{e}}, \hat{\mathbf{e}}) - (T\mathbf{x}, \mathbf{x})| &= |(T\hat{\mathbf{e}}, \hat{\mathbf{e}}) - (T\hat{\mathbf{x}}, \hat{\mathbf{x}}) - (T\hat{\mathbf{y}}, \hat{\mathbf{y}})| \\ &\leq (1 - \|\hat{\mathbf{x}}\|_{L_2(\Omega)}^2)(T\hat{\mathbf{e}}, \hat{\mathbf{e}}) + (T\hat{\mathbf{y}}, \hat{\mathbf{y}}). \end{aligned}$$

Combining (5.13), (5.14) and (5.18), we find

$$|(T\hat{\mathbf{e}}, \hat{\mathbf{e}}) - (T\mathbf{x}, \mathbf{x})| \leq C_\epsilon h^{4-\epsilon},$$

which together with (5.16)–(5.17) implies

$$(5.19) \quad |\mu - \mu_{h,\ell}| \leq C_\epsilon h^{2-\epsilon}.$$

The estimate (5.8) follows from (5.12), (5.15) and (5.19).

Recall that the gap $\hat{\delta}(M, N)$ between two subspaces M and N of a normed linear space $(X, \|\cdot\|_X)$ is defined by (cf. [39])

$$\hat{\delta}(M, N) = \max(\delta(M, N), \delta(N, M)),$$

where

$$\delta(M, N) = \sup_{\substack{x \in M \\ \|x\|_X=1}} \inf_{y \in N} \|x - y\|_X.$$

Therefore the statements about the gap between V_λ and $V_{h,\lambda}$ follow immediately from (5.9) and (5.10). \square

Remark 5.2. The compactness of the solution operator T and the existence of the uniform estimates (5.6) and (5.7) greatly simplify the analysis of the Maxwell eigenproblem. These ingredients are absent from Maxwell spectral approximations based on the full curl-curl problem and hence their justifications are much more involved [10, 11, 8, 20, 21, 9, 19, 12].

6. Numerical Results. In this section, we report a series of numerical examples that corroborate our theoretical results. Besides the L_2 error $\|\hat{\mathbf{u}} - \hat{\mathbf{u}}_h\|_{L_2(\Omega)}$ and the energy error $\|\hat{\mathbf{u}} - \hat{\mathbf{u}}_h\|_h$, we also report the errors measured in the semi-norm $|\cdot|_{\text{curl}}$ defined by

$$|\mathbf{v}|_{\text{curl}} = \|\nabla_h \times \mathbf{v}\|_{L_2(\Omega)}.$$

6.1. Source Problems. We first demonstrate the performance of our scheme for the source problems where we take α to be $-k^2$ for all the computations in this subsection.

In the first experiment, we check the convergence behavior of our numerical scheme (3.1) on the square $(0, 0.5)^2$ with uniform meshes, where the exact solution is

$$(6.1) \quad \hat{\mathbf{u}} = [y(y - 0.5) \sin(ky), x(x - 0.5) \cos(kx)]$$

TABLE 6.1

Convergence of the scheme (3.1) on the square $(0, 0.5)^2$ with uniform meshes and exact solution $\hat{\mathbf{u}}$ given by (6.1)

h	$\frac{\ \hat{\mathbf{u}} - \hat{\mathbf{u}}_h\ _{L_2(\Omega)}}{\ \hat{\mathbf{u}}\ _{L_2(\Omega)}}$	order	$\frac{\ \hat{\mathbf{u}} - \hat{\mathbf{u}}_h\ _h}{\ \hat{\mathbf{u}}\ _h}$	order	$\frac{ \hat{\mathbf{u}} - \hat{\mathbf{u}}_h _{\text{curl}}}{ \hat{\mathbf{u}} _{\text{curl}}}$	order
$k = 0$						
1/10	8.11E-02	–	3.61E-01	–	1.70E-01	–
1/20	1.85E-02	2.13	1.76E-01	1.04	8.52E-02	0.99
1/40	4.43E-03	2.07	8.69E-02	1.02	4.25E-02	1.00
1/80	1.08E-03	2.03	4.30E-02	1.01	2.12E-02	1.00
1/160	2.68E-04	2.02	2.15E-02	1.00	1.06E-02	1.00
$k = 1$						
1/10	7.19E-02	–	3.45E-01	–	1.64E-01	–
1/20	1.65E-02	2.12	1.68E-01	1.04	8.14E-02	1.01
1/40	3.93E-03	2.07	8.32E-02	1.02	4.09E-02	1.00
1/80	9.60E-04	2.04	4.13E-02	1.01	2.05E-02	1.00
1/160	2.37E-04	2.01	2.06E-02	1.01	1.02E-02	1.00
$k = 10$						
1/10	4.82E-01	–	9.11E-01	–	4.33E-01	–
1/20	7.62E-02	2.66	3.65E-01	1.32	2.00E-01	1.12
1/40	1.70E-02	2.17	1.75E-01	1.06	9.89E-02	1.02
1/80	4.07E-03	2.06	8.63E-02	1.02	4.93E-02	1.01
1/160	9.99E-04	2.03	4.32E-02	1.00	2.47E-02	1.00

for $k = 0, 1$ and 10 . The results tabulated in Table 6.1 confirm the estimates (4.24) and (4.25). Observe also that, as expected, finer meshes are needed for computing satisfactory approximate solutions when the wave numbers become larger.

Since our numerical scheme is designed for solving the reduced curl-curl problem (1.5), its performance should not be affected by the addition to the right-hand side of a gradient term ∇G where $G \in H_0^1(\Omega)$. This is demonstrated by our second experiment. Here we add ∇G with

$$(6.2) \quad G(x, y) = xy(x - 0.5)(y - 0.5) \sin(x + y)$$

to the right-hand side

$$(6.3) \quad \mathbf{f} = \nabla \times (\nabla \times \hat{\mathbf{u}}) - \hat{\mathbf{u}}$$

of the problem in the first experiment (with $k = 1$) and compare their performance. The results are reported in Table 6.2.

Note that the scheme (3.1) differs from the one proposed in [15] because of the last two stabilizing terms. In the third experiment, we show that these two terms are necessary for the convergence of our method due to the lack of the continuity of the approximating functions across the element interfaces. See Table 6.3 for the numerical evidence.

The last example in this subsection demonstrates the convergence behavior of our scheme on the L -shaped domain $(-0.5, 0.5)^2 \setminus [0, 0.5]^2$ with corners $(0.5, 0)$, $(0, 0)$, $(0, 0.5)$, $(-0.5, 0.5)$, $(-0.5, -0.5)$ and $(0.5, -0.5)$. We take $k = 1$ and the exact solution

TABLE 6.2

Comparison of the performance of the scheme (3.1) on $(0, 0.5)^2$ with $k = 1$ for the right-hand sides \mathbf{f} and $\mathbf{f} + \nabla G$, where \mathbf{f} is given by (6.3) and G is given by (6.2)

h	rhs = \mathbf{f}	rhs = $\mathbf{f} + \nabla G$
$\ \dot{\mathbf{u}} - \dot{\mathbf{u}}_h\ _{L_2(\Omega)}$		
1/10	0.001643092661	0.001643020363
1/20	0.000376001607	0.000375992406
1/40	0.000089765112	0.000089764564
1/80	0.000021922450	0.000021922530
1/160	0.000005415103	0.000005417204
$\ \dot{\mathbf{u}} - \dot{\mathbf{u}}_h\ _h$		
1/10	0.051109229470	0.051109792169
1/20	0.024919760832	0.024919911664
1/40	0.012296880765	0.012296952986
1/80	0.006107323566	0.006107362343
1/160	0.003043400530	0.003043424464

TABLE 6.3

Loss of convergence without the last two stabilizing terms in the scheme (3.1). The computation is on the square $(0, 0.5)^2$ with uniform meshes and the exact solution $\dot{\mathbf{u}}$ is given by (6.1) with $k = 1$

h	$\frac{\ \dot{\mathbf{u}} - \dot{\mathbf{u}}_h\ _{L_2(\Omega)}}{\ \dot{\mathbf{u}}\ _{L_2(\Omega)}}$	order	$\frac{ \dot{\mathbf{u}} - \dot{\mathbf{u}}_h _{\text{curl}}}{ \dot{\mathbf{u}} _{\text{curl}}}$	order
1/10	5.17E-00	—	2.61E-01	—
1/20	4.91E-00	0.22	1.71E-01	0.61
1/40	4.78E-00	0.04	1.33E-01	0.37
1/80	4.69E-00	0.03	1.20E-01	0.15

is chosen to be

$$(6.4) \quad \dot{\mathbf{u}} = \nabla \times \left(r^{2/3} \cos \left(\frac{2}{3}\theta - \frac{\pi}{3} \right) \phi(r/0.5) \right),$$

where (r, θ) are the polar coordinates at the origin and the cut-off function is given by

$$\phi(r) = \begin{cases} 1 & r \leq 0.25 \\ -16(r - 0.75)^3 \\ \quad \times [5 + 15(r - 0.75) + 12(r - 0.75)^2] & 0.25 \leq r \leq 0.75 \\ 0 & r \geq 0.75 \end{cases}$$

The meshes are graded around the re-entrant corner with the grading parameter equal to $1/3$. The results are tabulated in Table 6.4 and they agree with the estimates (4.24) and (4.25).

6.2. Maxwell Eigenproblem. In this subsection, we report the numerical results for the Maxwell eigenproblem.

In the first experiment, the computation is carried out on the square domain $(0, \pi)^2$ with uniform meshes. In this case, the exact eigenvalues are $r^2 + s^2$, $r, s = 0, 1, 2, 3, 4, \dots$ with $r^2 + s^2 \neq 0$. For instance, the first ten eigenvalues are 1, 1, 2, 4, 4,

TABLE 6.4

Convergence of the scheme (3.1) with graded meshes on the L -shaped domain $(-0.5, 0.5)^2 \setminus [0, 0.5]^2$ with $k = 1$ and exact solution $\hat{\mathbf{u}}$ given by (6.4)

h	$\frac{\ \hat{\mathbf{u}} - \hat{\mathbf{u}}_h\ _{L_2(\Omega)}}{\ \hat{\mathbf{u}}\ _{L_2(\Omega)}}$	order	$\frac{\ \hat{\mathbf{u}} - \hat{\mathbf{u}}_h\ _h}{\ \hat{\mathbf{u}}\ _h}$	order	$\frac{\ \hat{\mathbf{u}} - \hat{\mathbf{u}}_h\ _{\text{curl}}}{\ \hat{\mathbf{u}}\ _{\text{curl}}}$	order
1/4	1.77E+02	-	2.20E+01	-	9.13E-00	-
1/8	4.22E+01	2.07	7.76E-00	1.50	4.52E-00	1.02
1/16	3.42E-00	3.63	2.24E-00	1.79	7.86E-01	2.52
1/32	7.10E-01	2.27	1.09E-00	1.04	4.27E-01	0.88
1/64	1.77E-01	2.01	5.53E-01	0.98	2.34E-01	0.86

5, 5, 8, 9, 9. In Figure 6.1, we plot the first twenty numerical eigenvalues versus the parameter $n = \pi/h$. The symbol “o” on the right side denotes the exact eigenvalue, and “(2)” indicates that the multiplicity of the eigenvalue is 2. All the numerical approximations converge at a second order rate. To save space, only the first five numerical eigenvalues are included in Table 6.5. Furthermore, there is no spurious eigenvalue in our results.

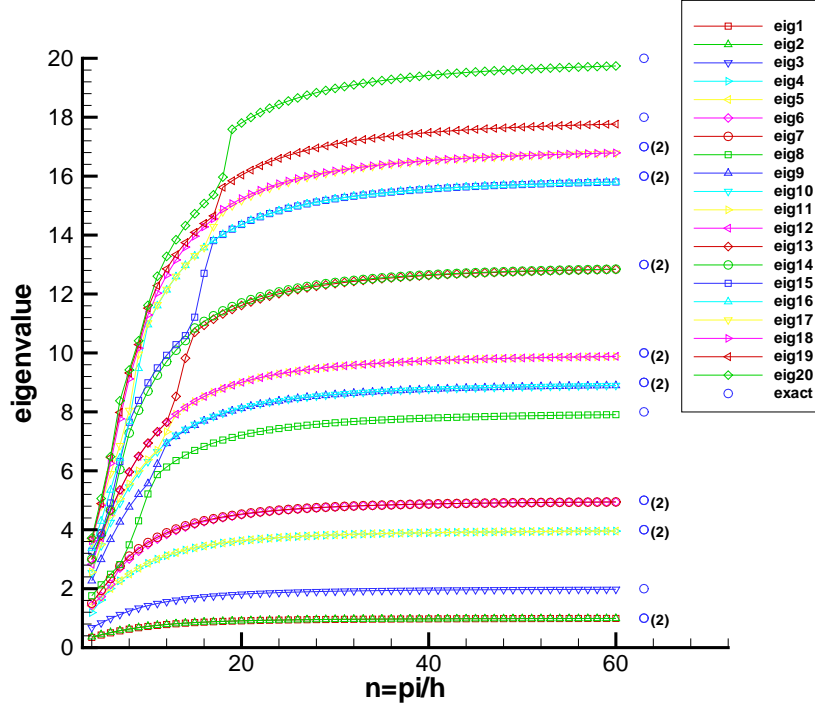


FIG. 6.1. First twenty numerical eigenvalues versus $n = \pi/h$ for Maxwell operator on $(0, \pi)^2$

In the second experiment, we compute the eigenvalues for the L -shaped domain $(-0.5, 0.5)^2 \setminus [0, 0.5]^2$ and the meshes are graded around the re-entrant corner with the grading parameter equal to $1/3$. Table 6.6 contains the first five numerical eigenvalues,

TABLE 6.5

Convergence of the first five numerical eigenvalues on $(0, \pi)^2$, with $\gamma =$ order of convergence and $\lambda =$ exact eigenvalue

h	1 st	γ	2 nd	γ	3 rd	γ	4 th	γ	5 th	γ
$\pi/4$	0.347	-	0.367	-	0.673	-	1.188	-	1.409	-
$\pi/8$	0.626	0.80	0.634	0.79	1.239	0.80	2.495	0.90	2.496	0.78
$\pi/16$	0.863	1.45	0.866	1.45	1.723	1.46	3.446	1.44	3.446	1.44
$\pi/32$	0.962	1.84	0.963	1.84	1.923	1.84	3.845	1.83	3.845	1.83
$\pi/64$	0.990	1.97	0.990	1.97	1.980	1.97	3.960	1.96	3.960	1.96
$\pi/128$	0.998	1.97	0.998	2.00	1.995	2.01	3.990	1.99	3.990	1.99
λ	1		1		2		4		4	

which show second order convergence of our method. This agrees with our analysis. In this case the “exact eigenvalues” are derived from numerical results of Dauge, which can be found at <http://perso.univ-rennes1.fr/monique.dauge/core/index.html>. In Figure 6.2, we plot the first ten numerical eigenvalues versus the parameter $n = 1/2h$. Again, there is no spurious eigenvalue.

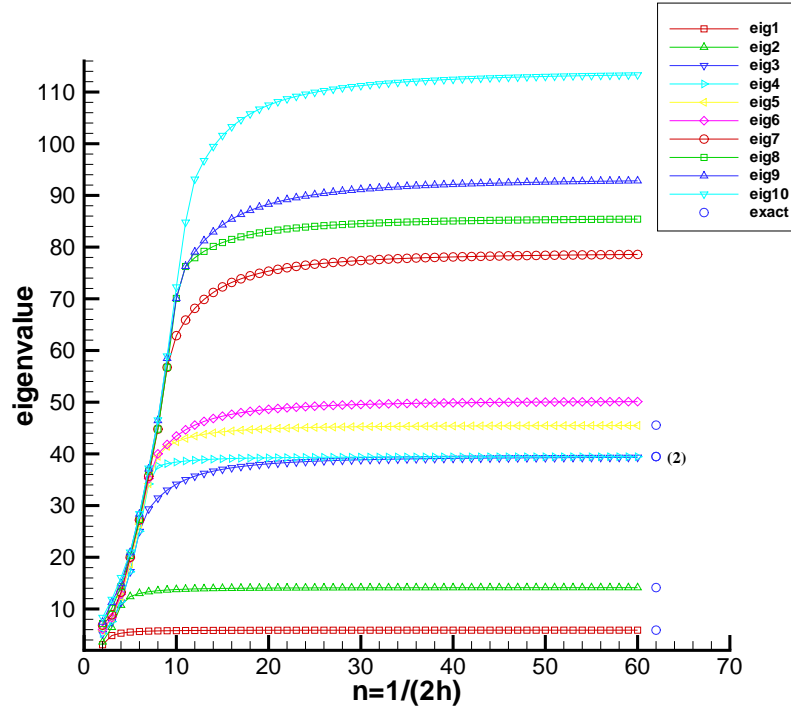


FIG. 6.2. First ten numerical eigenvalues versus $n = 1/2h$ for the Maxwell operator on $(-0.5, 0.5)^2 \setminus [0, 0.5]^2$.

7. Concluding Remarks. We have analyzed an interior penalty method with weak over-penalization for the two-dimensional reduced curl-curl problem and the

TABLE 6.6

Convergence of the first five numerical eigenvalues on the L-shaped domain $(-0.5, 0.5)^2 \setminus [0, 0.5]^2$ with graded meshes, $\mu = 1/3$ at the re-entrant corner, $\gamma =$ order of convergence and $\lambda =$ exact eigenvalue

h	1 st	γ	2 nd	γ	3 rd	γ	4 th	γ	5 th	γ
1/4	2.626	-	3.442	-	5.087	-	5.174	-	5.656	-
1/8	4.747	1.50	9.963	1.36	11.233	0.28	11.316	0.28	12.226	0.26
1/16	5.574	1.81	13.102	2.01	30.631	1.67	36.653	3.32	38.548	2.25
1/32	5.819	1.98	13.875	1.98	37.008	1.84	38.822	2.11	44.091	2.26
1/64	5.882	2.02	14.071	2.00	38.847	1.97	39.319	2.04	45.204	2.05
λ	5.902		14.136		39.478		39.478		45.558	

Maxwell eigenproblem. This method is inconsistent but stable without the need of tuning any penalty parameter. We expect other discontinuous Galerkin schemes will be discovered along this line. An advantage of this new approach is that fast solvers for the resulting schemes can be constructed using techniques developed for elliptic problems.

The scheme developed in this paper can be extended to three dimensional domains. In the case of uniform meshes, our two dimensional analysis can be easily generalized to establish convergence at a sub-optimal rate. It is likely that quasi-optimal convergence rates can be recovered using graded meshes, which is of course much more complicated in three dimensions [2, 3]. This will be one of the topics of our ongoing research.

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