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2 Title: A Nonlinear Transform for the Diagonalization of the Bernoulli-Laplace
3 Diffusion Model and Orthogonal Polynomials

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6 **Abstract.** The Bernoulli-Laplace model describes a diffusion process of two types of particles
7 between two urns. To analyze the finite-size dynamics of this process, and for other constructive
8 results we diagonalize the corresponding transition matrix and calculate explicitly closed-form ex-
9 pressions for all eigenvalues and eigenvectors of the Markov transition matrix T_{BL} . This is done by
10 a new method based on mapping the eigenproblem for T_{BL} to the associated problem for a linear
11 partial differential operator L_{BL} acting on the vector space of homogeneous polynomials in three
12 indeterminates. The method is applicable to other Two Urns models and is relatively easy to use
13 compared to previous methods based on orthogonal polynomials or group representations.

14 **Key word.** Urn models Diffusion Eigenvalues Generating Functions

15 **1. Introduction.** The Bernoulli-Laplace (BL) model arise from diffusion theory

16 and is related to the shuffling of cards [8]. Symmetries of the permutation group S_N
17 appear naturally in this model and other random walks on groups. Previous solutions
18 of this model have appeared in Diaconis and Shashahani [7] and in the works of
19 Karlin and MacGregor [10]. Group representations are used explicitly in the first; the
20 derivation of a non-standard inner product or equivalently a measure for orthogonal
21 polynomials which are related to the eigensolutions of the BL model appears in the
22 second.

23 In this paper, we give a third way for deriving exact solutions of all the eigen-
24 vectors of the BL model, through a nonlinear transform that triangularizes and then
25 diagonalizes the transition matrix T_{BL} . In brief, our method associates a specific
26 linear partial differential operator (LPDO) L_{BL} that acts on the vector space of ho-
27 mogeneous polynomials, G , to the matrix T_{BL} . The L_{BL} inherits the symmetries of
28 the BL model; it encodes the tri-diagonal singly-stochastic (column sums are equal to
29 1) and anti-symmetric structure of T_{BL} . The components of the (right) eigenvectors
30 (in view of the equal column sums of T_{BL}) of T_{BL} is encoded in the coefficients of
31 the homogeneous polynomial G . A classical theory for the symmetries of such LP-
32 DOs have been formulated in terms of the Lie algebra of symmetry operators K that
33 commutes with L_{BL} (cf. [13]).

34 It turns out and we exploit in our method, that the symmetries of L_{BL} appear
35 in the form of suitable linear and nonlinear transformations P on the independent
36 variables x, y , etc. or indeterminates of G . The ease of use of this method resides
37 in the transparent or explicit way to find these transformations P . Our algorithm is
38 completed by associating the transformed LPDO, L'_{BL} , back to what turns out to be
39 a triangular matrix T'_{BL} ; in other words, the transformation P for L_{BL} encodes a
40 similarity transformation that triangularizes T_{BL} , i.e., $PT_{BL}P^{-1} = T'_{BL}$, which is
41 then solved directly for its eigenvalues and right eigenvectors.

42 Here, we give a summary of the Urn models to which the BL model is related as
43 an extension. The Ehrenfest model and the Polya Urn models are two of the early

44 solvable models in the literature [9]. They appear as two of the exactly solved cases
 45 in Friedman’s formulation of Urn models where precisely one urn and balls of two
 46 colors are drawn and replaced with additions [9]. A dual formulation of Friedman’s
 47 Urn models was introduced in a series of recent papers [16, 14]: instead of balls of two
 48 colors and one urn, the dual formulation uses two urns and one-colored balls. The
 49 latter is more convenient for modelling of certain network science models [3], such
 50 as the Voter model where two balls are drawn and returned to the two urns with
 51 prescribed probabilities that depend on the order in which they are drawn. This is
 52 because many of the network science models are irreversible Markov chain models
 53 [4] which have absorbing states. Their transition matrices T , unlike T_{BL} for the BL
 54 model, are not symmetric, in an essential sense, that is, there are no non-standard
 55 inner products for R^N in which these matrices T have a symmetric form.

56 Using a new method based on diagonalization of transition matrices [15], we solved
 57 exactly the eigenvectors of several well-known models, including the Ehrenfest model,
 58 the Voter model [18, 12, 11, 5, 17, 2], the Moran model for genetic drift, and the
 59 Naming game models [19], [20]. Most of these models are irreversible Markov models
 60 with absorbing states, and have essentially non-symmetric transition matrices in the
 61 sense just mentioned. The BL model however, is based on two urns and balls of two
 62 colors. Thus, it is not strictly in the class of Two Urns models to which we recently
 63 applied our method. In modifying this method so that it applies to the BL model,
 64 we will have shown that the new method is not only easy to use but also flexible in
 65 extension to new problems.

66 One of the main points here is the technical simplicity of uncovering the symme-
 67 tries of the above LPDOs within our method, through the explicit appearance of the
 68 expressions $u = f(x, y)$ in the coefficients of L . We give here the LPDO L_V for the
 69 Voter model [16, 14] to indicate what we mean: first the propagation equation for the
 70 transition matrix T_V is given by

$$71 \quad (1.1) \quad a_j^{(m+1)} = p_{j-1}a_{j-1}^{(m)} + (1 - 2p_j)a_j^{(m)} + p_{j+1}a_{j+1}^{(m)}$$

$$72 \quad (1.2) \quad p_j = \frac{j(N-j)}{N(N-1)};$$

73 the eigen-problem for the associated LPDO $L_V = (x-y)^2 G_{xy}$ acting on the homo-
 74 geneous polynomial $G(x, y) = \sum_j c_j x^j y^{N-j}$ (which encodes the components c_j of the
 75 right-eigenvector of T_V) is given

$$76 \quad (1.3) \quad (x-y)^2 G_{xy} = N(N-1)(\lambda-1)G,$$

77 which clearly suggests the transformation $u = x - y$, $v = y$. Indeed this triangularized
 78 and diagonalized the Voter model and led to its complete solution.

79 Contrast this ease of use with the fact that triangularization and diagonalization
 80 of a given transition matrix of size N has computational complexity $O(N^3)$. In other
 81 words, exact integration of the Two Urns models via diagonalization of transition
 82 matrices are nontrivial problems, that are difficult to solve but once known, the so-
 83 lutions are easy to verify. [16, 14] provides a simple method to find such explicit
 84 diagonalization and hence all eigenvectors for a class of transition matrices from the
 85 Two Urns models, even when their transition matrices are essentially non-symmetric.
 86 Note that the eigenproblem and diagonalization of symmetric matrices have a lower
 87 computational complexity.

88 We aim here to highlight this method's ease of use, relative to the group rep-
89 resentation method and the method of orthogonal polynomials. Moreover, the BL
90 model differs significantly from the original Two Urns subclass of models for which
91 our method was initially formulated. Thus, we also aim to show that, with the spe-
92 cific introduction of a nonlinear change of independent variables, this method can be
93 applied to more complex models than the original class of models. Since the BL tran-
94 sition matrix T_{BL} is from a reversible Markov chain with a stationary distribution [4],
95 it is non-symmetric singly-stochastic only in a trivial sense. In other words, there ex-
96 ists (a difficult to find) non-standard inner product for R^N , in which T_{BL} is symmetric
97 and hence doubly-stochastic. Karlin and McGregor [10] using their powerful integral
98 representation method of finding an explicit way to symmetrize T by introducing a
99 non-standard inner product (or orthogonal measure) into the problem, have related
100 the right eigenvectors of the transition matrix T_{BL} of the BL model to the orthogonal
101 polynomials called the dual Hahn polynomials. A third aim of this paper is there-
102 fore to re-derive from the diagonalization of T_{BL} , this non-standard inner product in
103 which the dual Hahn polynomials are an orthogonal polynomial system. Note that
104 this non-standard inner product, once found, yields a symmetric version of T_{BL} which
105 is an example of a Jacobi operator that arise in the the classical moments problem
106 [1], and is related to orthogonal polynomials via the Riemann-Hilbert method [6].

107 The paper is organized as follows: section 3 concerns the calculation of the right
108 eigenvectors and eigenvalues of T in closed form by method introduced in [16, 15,
109 14]; in view of the fact that these right eigenvectors are not orthogonal in the usual
110 Euclidean inner-product, section 4 concerns the transformations needed to calculate
111 the orthonormal system of left eigenvectors of T_{BL} , and also the derivation of the non-
112 standard inner product in which the right eigenvectors are now orthogonal; section 5
113 concerns the elementary proofs, based only on the eigenvectors and eigenvalues of T ,
114 for the tight upper and lower bounds for times to stationarity in the BL model [7] ;
115 section 6 concerns a numerically exact evaluation of the expression for the TV norm
116 in these bounds using the eigenvectors and eigenvalues of T_{BL} directly, hence slightly
117 sharper estimates for the mixing times of the BL model.

118 Beyond the balanced special case of the BL problem treated in detail in this paper,
119 the same generating function method can be used to prove similar tight bounds for
120 mixing times in the other cases.

121 **2. Transition matrix of the BL model.** Let the transition matrix T_{BL} be
122 defined so that $(T_{BL})_{ij} = Pr\{n_w(m+1) = i | n_w(m) = j\}$, so that the sum of each
123 column is 1. In the general BL model for balls of two colors and two urns, $N_1, N_2,$
124 N_w, N_b are fixed parameters satisfying the constraints

$$125 \quad (2.1) \quad N_w + N_b = N = N_1 + N_2$$

126 where N equals total number of balls in the model. For $i = 0, \dots, N_w \leq N_1$, (where by
127 abuse of notation i stands for both the row label of transposed matrix T_{BL}^t and the
128 number of white balls in urn 1, n_w), the transition probabilities are explicitly given
129 by

$$130 \quad (2.2) \quad p_i = Pr\{n_w(t+1) = i+1 | n_w(t) = i\} = \frac{(N_1 - i)(N_w - i)}{N_1 N_2}$$

$$131 \quad (2.3) \quad q_i = Pr\{n_w(t+1) = i-1 | n_w(t) = i\} = \frac{i(N_b - (N_1 - i))}{N_1 N_2}$$

$$132 \quad (2.4) \quad r_i = Pr\{n_w(t+1) = i | n_w(t) = i\} = 1 - q_i - p_i.$$

134 **3. Diagonalization - Right eigenvectors of the general BL model.** In
 135 [16, 14], we developed an explicit method for exactly integrating or solving a 5-
 136 parameters subclass of a class of Two Urns models which is parametrized by six real
 137 parameters. Our method is based on a relationship between certain banded stochastic
 138 matrices T (such as tridiagonal and pentadiagonal non-symmetric transition matrices
 139 of markov chain models) and the LPDOs acting on the vector space of homogeneous
 140 polynomials, $G(x, y)$ of finite order in two indeterminates. The symmetries of the
 141 LPDO, L , associated with a given non-symmetric singly stochastic matrix from this
 142 solvable subclass of the Two Urns models, are identified and used explicitly to trans-
 143 form from the original indeterminates (independent variables x, y say) to suitable new
 144 variables (such as $u = f(x, y), v = g(x, y)$). In the new variables u, v , the transformed
 145 LPDO, L' , acts on the (again homogeneous of same order as $G(x, y)$) polynomial
 146 $H(u, v)$. We have shown in [16, 14] that the transformed eigen-problem

$$147 \quad (3.1) \quad L'[H(u, v)] = \lambda(N)H(u, v)$$

148 for a well-defined subclass of such Two Urns problems is equivalent (via the inverse
 149 of the original relationship between banded matrix and LPDO) to the eigen-problem
 150 for a triangular matrix, which can then be solved explicitly for both right and left
 151 eigenvectors. In other words, at the end of this brief summary, the symmetries of
 152 LPDO L inherited from the original banded stochastic matrix T , generate an explicit
 153 similarity transformation, P , such that

$$154 \quad (3.2) \quad STS^{-1} = D$$

155 where D is diagonal, and S contains the eigenvectors of T .

156 This method can be formalized as an Algorithm as follows: Given the input of a
 157 singly stochastic transition matrix T of size $N + 1$,

158 (I) Choose a suitable homogeneous polynomial of finite degree N , G that has the
 159 components c_i of a right eigenvector of T as coefficients of the monomials $x^i y^j z^k$; part
 160 of this choice is the number of indeterminates in G . For example, the Voter model of
 161 size N (number of balls) with a transition matrix $T_V(N)$ which is a $N + 1$ by $N + 1$
 162 real matrix, requires a homogeneous polynomial G_V of degree N in the indeterminates
 163 x, y because there are two urns.

164 (II) Associate the recursion implicit in given Markov matrix T to a LPDO, L
 165 which acts on the homogeneous polynomial G ; the basic elements of this association
 166 scheme are the standard linear differential operators for increasing, decreasing and not
 167 changing the numbers of balls in each urn (which correspond in the example below to
 168 the probabilities p, q, r prescribed by the transition matrix), and a set of multiplication
 169 type linear operators that correspond to shifts.

170 (III) A transformation to new independent variables, (for instance, $u = f(x, y, z)$,
 171 $v = g(x, y, z)$, $w = h(x, y, z)$) is chosen to satisfy two conditions:

172 (A) the transformed polynomial

$$173 \quad H(u, v, w) = H(f(x, y, z), g(x, y, z), h(x, y, z)) = G(x, y, z)$$

174 is a homogeneous polynomial of the same finite degree as G ;

175 (B) $u = f(x, y, z)$, $v = g(x, y, z)$, $w = h(x, y, z)$ is a transformation based on the
 176 symmetries of L (cf. [13]), that is, the combinations $f(x, y, z)$, $g(x, y, z)$, $h(x, y, z)$ of
 177 the original variables x, y, z appear naturally in the coefficients of the LPDO, L .

178 These conditions (A) and (B) are clearly not sufficient to ensure the transformed
 179 LPDO eigenproblem

$$180 \quad (3.3) \quad L'[H(u, v, w)] = \lambda(N)H(u, v, w)$$

181 is associated with a similar triangular matrix T' which explicitly yields all its eigen-
 182 vectors b_i . That they are sufficient has to be proved either in each problem to which
 183 we apply the Algorithm, or for a class of models as in the case of the Two Urns
 184 models.

185 (IV) Using the transformation in step (III), derive the corresponding transformed
 186 LPDO, L' that acts on the transformed polynomial $H(u, v, w)$.

187 (V) Without explicitly calculating the transformed matrix T' which is associ-
 188 ated with the transformed LPDO, L' in step (IV), check that the transformed eigen-
 189 problem for L' is indeed a recursion system for the transformed eigenvectors b_i that
 190 can be solved explicitly, i.e., it is equivalent to a triangular linear system of equations.
 191 Solve for the eigenvalues and then the eigenvector components b_i , and if required
 192 transform back to the original components c_i . These are the main outputs of the
 193 Algorithm.

194 (VI) Use the eigenvectors in step (V) to diagonalize the original matrix if neces-
 195 sary. This is the end of the ALgorithm.

196 Now we apply the Algorithm to the BL model. Given the transition matrix of
 197 the BL model (cf. section 2), it will be obvious that three independent variables
 198 (instead of the two before) should be used to formulate the BL problem. In step
 199 (I) of the Algorithm, we adopt the ansatz that the LPDO, L_{BL} , associated with
 200 the above N by N matrix T_{BL} , now acts on a homogeneous polynomial $G(x, y, z)$ in
 201 three indeterminates, x, y, z . We encode the entries $c_k(i)$, $i = 0, \dots, N_w$ of the k -th
 202 eigenvector of the transition matrix for the BL model as follows:

$$203 \quad (3.4) \quad G^{(k)}(x, y, z) = \sum_{i=0}^{N_w} c_k(i) x^i y^{N_1-i} z^{N_w-i}.$$

204 where i = number of white balls in urn 1 (also denoted n_w). The choice of three
 205 independent variables to encode the components of an eigenvector of T_{BL} in the
 206 homogeneous polynomial $G(x, y, x)$ is now made obvious by this explicit expression
 207 for G .

208 In step (II) of the Algorithm, we derive from the original eigen-problem for tran-
 209 sition matrix T_{BL} , an LPDO, L_{BL} , that acts on $G^{(k)}$. Towards that aim, we note, in
 210 particular, the entries for p_i and q_i in T_{BL} correspond respectively to the following
 211 linear differential operators with coefficients that are monomials in $x, y,$ and z ,

$$212 \quad (3.5) \quad L_p = \frac{yzG_{yz}^{(k)}}{N_1N_2}$$

$$213 \quad (3.6) \quad L_q = \frac{N_b x}{N_1 N_2} G_x^{(k)} - \frac{xy}{N_1 N_2} G_{xy}^{(k)},$$

214 where $G_{yz}^{(k)} = \frac{\partial^2}{\partial y \partial z} G^{(k)}$ for example. In addition, it is part of the association scheme
 215 that multiplication in the LPDO (cf. [13]) by the coefficient $\frac{x}{yz}$ (resp. $\frac{yz}{x}$) represents
 216 down (resp. up) shifts in the index i within the discrete recursion equations of the
 217 original eigen-problem for matrix T_{BL} . The L_{BL} associated with the eigen-problem
 218 of the tridiagonal Markov matrix T_{BL} is given by:

$$219 \quad (3.7) \quad L_{BL}[G^{(k)}] = N_1 N_2 (\lambda_k - 1) G^{(k)}$$

$$220 \quad (3.8) \quad L_{BL}[G^{(k)}] \equiv (x - yz) G_{yz}^{(k)} + y(x - yz) G_{xy}^{(k)} - N_b (x - yz) G_x^{(k)}.$$

221 In step(III) of the Algorithm, we note that the symmetries of L_{BL} with respect
 222 to transformations of its independent variables, is expressed in the factor $(x - yz)$ in
 223 its coefficients. This suggests the transformation to the new independent variables

$$224 \quad (3.9) \quad u = x - yz, y = y, z = z.$$

225 Since the transformed homogeneous polynomial is now given by

$$226 \quad (3.10) \quad H^{(k)}(u, y, z) = G^{(k)}(x(u, yz), y, z) = \sum_i b_i^k u^i y^{N_1 - i} z^{N_w - i}$$

227 in terms of the (new) components b_i^k of the k -th right eigenvector, this transformation
 228 clearly satisfies both necessary conditions (A) and (B) in step (III) of the Algorithm.

229 To prove that it is sufficient for our purpose of obtaining the eigenvalues and
 230 eigenvectors exactly and for diagonalization, we proceed by direct calculations.

231 In step (IV) using the following obvious identities for the transformation of partial
 232 derivatives

$$233 \quad (3.11) \quad \partial_x = \partial_u$$

$$234 \quad (3.12) \quad \partial_y = \partial_y - z \partial_u$$

$$235 \quad (3.13) \quad \partial_z = \partial_z - y \partial_u$$

$$236 \quad (3.14) \quad \partial_{xy} = \partial_{yu} - z \partial_u^2$$

$$237 \quad (3.15) \quad \partial_{yz} = \partial_{yz} - y \partial_{yu} - z \partial_{uz} + yz \partial_u^2 - \partial_u$$

238 the transformed LPDO, L'_{BL} in $H^{(k)}$, $k = 0, \dots, N_w$ is

$$239 \quad (3.16) \quad N_1 N_2 (\lambda_k - 1) H^{(k)} = L'_{BL}[H^{(k)}],$$

$$240 \quad (3.17) \quad L'_{BL}[H^{(k)}] = -N_b u \partial_u H^{(k)} + yu (\partial_{yu} - z \partial_u^2) H^{(k)}$$

$$241 \quad (3.18) \quad + u (\partial_{yz} - y \partial_{yu} - z \partial_{uz} + yz \partial_u^2 - \partial_u) H^{(k)}$$

$$242 \quad (3.19) \quad = u (\partial_{yz} - z \partial_{uz}) H^{(k)} - (N_b + 1) u \partial_u H^{(k)}$$

243 In step (V), by reversing the derivation of the original L_{BL} through the association
 244 scheme [13], this L'_{BL} in H is shown to be equivalent to the following triangular system
 245 for the (right) eigen-problem of the transformed matrix T'_{BL} :

$$246 \quad (3.20) \quad N_1 N_2 (\lambda_k - 1) b_i^k$$

$$247 \quad (3.21) \quad = (N_1 - i + 1) (N_w - i + 1) b_{i-1}^k - i (N_w - i) b_i^k - (N_b + 1) i b_i^k.$$

248 We have therefore verified the sufficiency of the transformation where $u = x - yz$
 249 for triangularizing (and later diagonalizing) T_{BL} . This triangular system implies the
 250 recursion

$$251 \quad (3.22) \quad b_i^k = \frac{(N_1 - i + 1) (N_w - i + 1) b_{i-1}^k}{N_1 N_2 (\lambda_k - 1) + i (N_w - i) + (N_b + 1) i}$$

252 which can be solved directly.

253 For nontrivial eigensolutions for $k = 0, \dots, N_w$, the denominator in b_i^k must vanish,
 254 yielding the following exact expressions for the eigenvalues,

$$255 \quad (3.23) \quad \lambda_k = 1 - \frac{k(1-k+N_w+N_b)}{N_1 N_2}$$

$$256 \quad (3.24) \quad = 1 - \frac{k(1-k+N)}{N_1 N_2}$$

$$257 \quad (3.25) \quad = 1 - \frac{k(N-k+1)}{N_1 N_2}$$

258 In the case $N_1 = N_w$,

$$259 \quad (3.26) \quad \lambda_0 = 1$$

$$260 \quad (3.27) \quad \lambda_1 = 1 - \frac{N}{N_1 N_2}.$$

261 The eigenvectors (in the transformed variables of H) are given explicitly by:

$$262 \quad (3.28) \quad b_i^k = \prod_{j=k+1}^i \frac{(N_1 - j + 1)(N_w - j + 1)}{N_1 N_2 (\lambda_k - 1) + j(N_w + N_b + 1 - j)}$$

$$263 \quad (3.29) \quad = \prod_{j=k+1}^i \frac{(j - N_1 - 1)(j - N_w - 1)}{-k(N + 1 - k) + j(N + 1 - j)}$$

$$264 \quad (3.30) \quad = \prod_{j=k+1}^i \frac{(j - N_1 - 1)(j - N_w - 1)}{(j - k)(j + k - N - 1)}$$

$$265 \quad (3.31) \quad = (-1)^{i-k} \frac{(k - N_1)_{i-k} (k - N_w)_{i-k}}{(i - k)! (2k - N)_{i-k}}.$$

266 Using these coefficients in the definition for H gives

$$267 \quad (3.32) \quad H^{(k)} = \sum_{i=k}^{N'} (-1)^{i-k} \frac{(k - N_1)_{i-k} (k - N_w)_{i-k}}{(i - k)! (2k - N)_{i-k}} u^i y^{N_1 - i} z^{N_w - i}$$

268 We summarize the consequences of the above steps of the Algorithm on the BL
 269 model in the following theorem:

270 **THEOREM 3.1.** *In the above Algorithm for the BL model, for any size N of the*
 271 *model, the LPDO, L'_{BL} , after the transformation (3.9) on the independent variables,*
 272 *is equivalent to a triangular linear system (3.21) which has (right) eigenvectors given*
 273 *by (3.31) and eigenvalues (3.25). The (right) eigenvectors of the original BL matrix*
 274 *T_{BL} are in turn given by (3.38).*

275 **3.1. Hypergeometric functions and dual Hahn polynomials.** The next to
 276 final step left in this part of the paper is step (VI) in the Algorithm, to invert the
 277 above similarity transformation to obtain explicitly the closed-form expressions for
 278 the original components of the right-eigenvectors $c_k(i)$ of T_{BL} . For this purpose, let
 279 $h^{(k)}(u) = H^{(k)}(u, 1, 1)$. Then,

$$280 \quad g^{(k)}(x) = G^{(k)}(x, 1, 1) = H^{(k)}(x - 1, 1, 1) = h^{(k)}(x - 1)$$

$$281 \quad = (x - 1)^k {}_2F_1(k - N_1, k - N_w; 2k - N; 1 - x).$$

282 Using the hypergeometric identity [],

283 (3.33)
$${}_2F_1(a, b; c; 1 - z) \propto {}_2F_1(a, b; a + b - c + 1; z)$$

284 and the fact that any multiple of an eigenvector remains an eigenvector, we take the
285 polynomial for the right eigenvector components to be

286 (3.34)
$$g^{(k)}(x) = (x - 1)^k {}_2F_1(k - N_1, k - N_w; N_2 - N_w + 1; x)$$

287 whose coefficients are the original components $c_k(i)$ of the k -th right eigenvector cor-
288 responding to λ_k prior to the transformation above. These expressions are equivalent
289 to the dual Hahn polynomials [10].

290 For the hypergeometric representation of the eigenvectors to be well defined, we
291 require $N_2 \geq N_w$. There is no loss in generality with this assumption, because we can
292 relabel $N_1 \leftrightarrow N_2$ and $N_w \leftrightarrow N_b$ so that the assumption holds. From the solution for
293 $g^{(k)}$, we expand in x^l ,

294 (3.35)
$$g^{(k)}(x) = \sum_n \binom{k}{n} (-1)^{k-n} x^n \sum_i \frac{(k - N_1)_i (k - N_w)_i}{(N_2 - N_w + 1)_i i!} x^i$$

295 (3.36)
$$= \sum_i \sum_n \binom{k}{n} (-1)^{k-n} \frac{(k - N_1)_i (k - N_w)_i}{(N_2 - N_w + 1)_i i!} x^{i+n}$$

296 (3.37)
$$= (-1)^k \sum_i \left[\sum_n \binom{k}{n} (-1)^n \frac{(k - N_1)_{i-n} (k - N_w)_{i-n}}{(N_2 - N_w + 1)_{i-n} (i - n)!} \right] x^i$$

297

298 to find the explicit form for the components of the k -th right eigenvectors:.

299 (3.38)
$$c_i^k = \sum_n \binom{k}{n} (-1)^n \frac{(k - N_1)_{i-n} (k - N_w)_{i-n}}{(N_2 - N_w + 1)_{i-n} (i - n)!}$$

300 Notice that the solution is the k th order backwards difference of the components of
301 the hypergeometric coefficients of g .

302 The above treatment of the eigen-problem by transforming via symmetries, the
303 independent variables of the associated LPDO, L_{BL} , is equivalent to a similarity
304 transformation of the transition matrix T_{BL} . Let $\mathbf{w} = \mathbf{P}\mathbf{v}$ for some transformation
305 matrix \mathbf{P} . Then, the eigen-problem for \mathbf{w} is given by $\mathbf{P}T_{BL}\mathbf{P}^{-1}\mathbf{w} = \lambda\mathbf{w}$. The
306 above calculations is equivalent to the matrix \mathbf{P} such that the new matrix $\mathbf{T}'_{BL} =$
307 $\mathbf{P}T_{BL}\mathbf{P}^{-1}$ is lower triangular. The last step in this section is to diagonalize T_{BL} .
308 We do this by diagonalizing the matrix triangular matrix $\mathbf{T}'_{BL} = \mathbf{W}\mathbf{\Lambda}\mathbf{W}^{-1}$. Here,
309 $\mathbf{\Lambda} = \text{diag}(\lambda_0, \dots, \lambda_N)$ and \mathbf{W} are the eigenvectors of \mathbf{T}'_{BL} . The components of these
310 eigenvectors are b_i corresponding to eigenvalue λ_k . Diagonalization of \mathbf{T}'_{BL} allows us
311 to explicitly diagonalize the original transition matrix as

312 (3.39)
$$\mathbf{T}_{BL} = \mathbf{P}^{-1}\mathbf{W}\mathbf{\Lambda}\mathbf{W}^{-1}\mathbf{P}.$$

313 Note that the matrix of eigenvectors is given by $\mathbf{P}^{-1}\mathbf{W}$.

314 **4. Symmetrizing transform, orthogonal measure and dual Hahn poly-**
315 **nomials.** For transition matrix T_{BL} , let Z be given by (where we drop the subscript
316 BL herein, i.e., $T = T_{BL}$)

317 (4.1)
$$Z_{ij} = \sqrt{\pi_j} T_{ij} \frac{1}{\sqrt{\pi_i}}.$$

318 Recall the detailed balance of T and its stationary distribution, given by $T_{ij}\pi_j = T_{ji}\pi_i$
 319 [4] follows from the reversibility and ergodicity of the BL model. Note that Z is the
 320 symmetric version of the transition matrix:

$$321 \quad (4.2) \quad Z_{ij} = \frac{1}{\sqrt{\pi_j}} \pi_j T_{ij} \frac{1}{\sqrt{\pi_i}}$$

$$322 \quad (4.3) \quad = \frac{1}{\sqrt{\pi_j}} T_{ji} \sqrt{\pi_i}$$

$$323 \quad (4.4) \quad = Z_{ji}.$$

325 Therefore, Z has an orthonormal set of left eigenvectors, w_k^T . Let W be a matrix
 326 whose columns are w_k . The spectral decomposition of Z by left eigenvectors is given
 327 by

$$328 \quad (4.5) \quad Z = W \Lambda W^T.$$

329 By the definition of Z_{ij} , the transformation from T to Z can be expressed as

$$330 \quad (4.6) \quad D^{-1} T D = Z,$$

331 where D is a diagonal matrix whose diagonal entries are $\sqrt{\pi_i}$. So, arbitrary powers
 332 of T is given by

$$333 \quad (4.7) \quad T^m = D W \Lambda^m (D^{-1} W)^T.$$

334 Defining a new transformation by S ,

$$335 \quad (4.8) \quad T^m = S \Lambda^m S^{-1}.$$

336 Since W has a specific normalization, we can equate S with DW after applying the
 337 appropriate normalization for S . That is, for diagonal matrix Δ , we take

$$338 \quad (4.9) \quad S \Delta = D W.$$

339 We can choose any normalization for the right eigenvectors given in S , and Δ will
 340 properly renormalize them. Here, we solve for Δ and W by appealing to the orthog-
 341 onality of W :

$$342 \quad (4.10) \quad W^T W = \Delta S^T D^{-2} S \Delta = I.$$

343 Therefore

$$344 \quad (4.11) \quad S^T D^{-2} S = \Delta^{-2}.$$

345 Computing the matrix multiplication on the left side yields the diagonal entries
 346 of Δ denoted by Δ_k given by

$$347 \quad (4.12) \quad \Delta_k^{-2} = \sum_{i=0}^N \frac{1}{\pi_i} c_k(i)^2$$

348 in terms of the right eigenvectors of T_{BL} . Now that we have Δ , we have the repre-
 349 sentations for both the left-eigenvectors $w_k(i)$ and right-eigenvectors $v_k(i)$ of Z given
 350 by

$$351 \quad (4.13) \quad w_k(i) = \frac{\Delta_k}{\sqrt{\pi_i}} c_k(i)$$

$$352 \quad (4.14) \quad v_k(i) = \frac{1}{\sqrt{\pi_i}} w_k(i) = \frac{\Delta_k}{\pi_i} c_k(i)$$

354 in terms of the right-eigenvectors $c_k(i)$ of the original T_{BL} that was obtained by our
 355 method in section 2.

356 From Eq. (4.11), we also have an explicit formula for S^{-1} given by

$$357 \quad (4.15) \quad S^{-1} = \Delta^2 S^T D^{-2}.$$

358 So, by Eq. (??), we have

$$359 \quad (4.16) \quad T^m = S \Lambda^m \Delta^2 S^T D^{-2}.$$

360 Computing the matrix multiplication gives the following spectral decomposition

$$361 \quad (4.17) \quad T_{ij}^{(m)} = \frac{1}{\pi_j} \sum_{k=0}^N \Delta_k^2 \lambda_k^m c_k(i) c_k(j).$$

362 as the explicit representation of $Pr\{n(m) = i \mid n(0) = j\}$ in the BL model.

363 Since $T^0 = I$, take $m = 0$ in Eq. (4.17) to find the stationary distribution of the
 364 BL model

$$365 \quad (4.18) \quad \pi_j \delta_{ij} = \sum_{k=0}^N \Delta_k^2 c_k(i) c_k(j).$$

366 Note that this is the orthogonality relation for the right-eigenvectors of T_{BL} with
 367 orthogonal measure Δ_k^2 given in Eq. (??). We have derived the orthogonal measure
 368 Δ_k^2 in which the dual Hahn are an orthogonal polynomial system [10].

369 We summarize these results on the derivation of a non-standard inner product or
 370 orthogonal measure in which the original transition matrix of the BL model becomes
 371 a symmetric real matrix and the (right) eigenvectors are the system of orthogonal
 372 dual Hahn polynomials:

373 **THEOREM 4.1.** *The orthogonal measure in (4.12) symmetrizes T_{BL} , is related to*
 374 *the (left) and (right) eigenvectors of T_{BL} by (4.13) and (4.14), and yields the spectral*
 375 *decomposition (4.17).*

376 **5. Bounds of mixing times - elementary proofs.** We will discuss first the

377 case $N_1 = N_2 = N/2$, for our method gave the eigenvalues of the BL model to be
 378 $\lambda_k = 1 - \frac{4k(N-k+1)}{N^2}$, $\lambda_1 = 1 - \frac{4}{N}$. A heuristic estimate of the number of switches q
 379 needed to mix the colors is thus,

$$380 \quad (5.1) \quad \left(1 - \frac{4}{N}\right)^q \simeq e^{-\frac{4q}{N}} = \frac{1}{N},$$

381 and therefore $q = \frac{1}{4}N \log N$ gives an idea of how many switches or time steps are
 382 needed until the variation distance between ρ_q and π is of order $O(1/N)$. The lower
 383 bound is obtained along the lines of Diaconis et al, that is, by an application of
 384 the Chebyshev's inequality. However, all estimates of the relevant mean values and
 385 variance needed to apply the Chebyshev's inequality are constructed explicitly from
 386 the properties of the eigenvalues and eigenvectors of the BL model. We will state the
 387 theorems below but since the proofs are similar to those in [7], we have provided the
 388 details in an appendix.

389 **THEOREM 5.1.** *For $m = \frac{1}{4}N \log N + (\frac{c}{2} - \frac{\log 2}{8})N$, for $c > 0$, there is a universal*
 390 *constant $A > 0$ such that $E_\pi [\|\rho_m(j; \cdot) - \pi\|_V] \leq Ae^{-2c}$.*

391 **THEOREM 5.2.** *If $m = \frac{1}{8}N \ln N - \frac{cN}{2}$, then $2\|\rho_m - \pi\|_V \geq 1 - e^{4c}$*

392 **6. Exact calculations of mixing times.** Given that we can calculate $\rho_m(i, j)$

393 exactly by Eq. (??), and the stationary distribution is given by

394 (6.1)
$$\pi_i = \frac{\binom{N_w}{i} \binom{N-N_w}{N_1-i}}{\binom{N}{N_1}},$$

395 the total variational distance can be exactly computed for all time steps by

396 (6.2)
$$\|\rho - \pi\|_V = \frac{1}{2} \sum_{i=0}^{N/2} \left| \pi_i \sum_{k=1}^{N/2} \lambda_k^m v_k(i) v_k(0) \right|$$

397 This solution is shown in Figure 1, with the upper bound given in Figure 2.

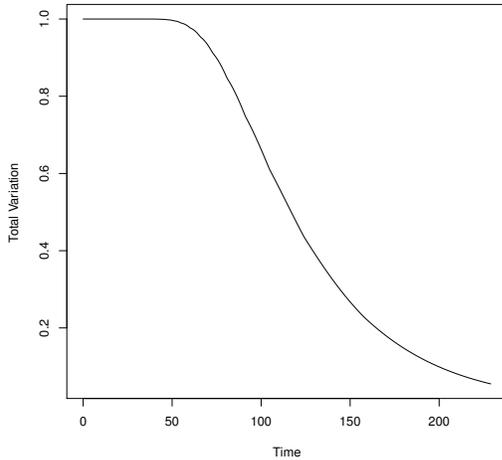


FIG. 1. Plot of the exact solution for the total variational distance for $N_1 = N_2 = N_w = N_b = 100$.

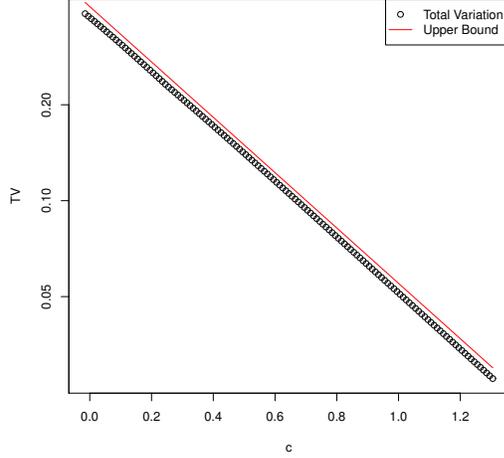


FIG. 2. Plot of the exact solution for the total variational distance, and the upper bound given in Theorem 5.1

398 **Appendix A. Proofs of Theorems 5.1 and 5.2.** The eigenvalues and eigen-
 399 vectors of the BL problem can be used to construct an upper bound on the variation
 400 distance between *the* probability distribution after m . As we proved in section 3, the
 401 right-eigenvectors of the original matrix T_{BL} are not orthogonal in the standard inner
 402 product of R^N but a different inner product weighted by Δ_k^2 can be used to derive
 403 orthogonality of a related system of right-eigenvectors v_j of $T = T_{BL}$.

404 **A.1. Upper bound.** Let $j = 0$ to define

$$405 \quad \rho_m(i; 0) = \pi_i + \pi_i \sum_{k=0}^{N/2} \lambda_k^m v_k(i) v_k(0)$$

406 where

$$407 \quad (\text{A.1}) \quad \Pr\{n_w(m) = i \mid n_w(0) = j\} = T_{ij}^m = \rho_m(j; i)$$

$$408 \quad (\text{A.2}) \quad = \sum_{k=0}^{N/2} \pi_i \lambda_k^m v_k(i) v_k(j)$$

$$409 \quad (\text{A.3}) \quad = \pi_i + \pi_i \sum_{k=1}^{N/2} \lambda_k^m v_k(i) v_k(j).$$

410 Then, $\rho_0(i; j) = \delta_{ij}$, and hence

$$411 \quad (\text{A.4}) \quad \sum_{k=1}^{N/2} v_k^2(i) = \frac{1}{\pi_i} - 1 < \frac{1}{\pi_i},$$

412 implies that

$$413 \quad (\text{A.5}) \quad \|\rho_m(i; 0) - \pi_i\|_V = \frac{1}{2} \sum_{i=0}^{N/2} \left| \pi_i \sum_{k=1}^{N/2} \lambda_k^m v_k(i) v_k(0) \right|$$

$$414 \quad (\text{A.6}) \quad \leq \frac{1}{2} \lambda_1^m \sum_{i=0}^{N/2} \pi_i \left| \sum_{k=1}^{N/2} v_k(i) v_k(0) \right|$$

$$415 \quad (\text{A.7}) \quad \leq \frac{1}{2} (N)^{-1/2} e^{-2c} \sum_{i=0}^{N/2} \pi_i \left(\sum_{k=1}^{N/2} v_k^2(i) \right)^{1/2} \left(\sum_{k=1}^{N/2} v_k^2(0) \right)^{1/2}$$

$$416 \quad (\text{A.8}) \quad \leq \frac{1}{2} (N \pi_0)^{-1/2} e^{-2c} \sum_{i=0}^{N/2} (\pi_i)^{1/2}$$

417 Using

$$418 \quad (\text{A.9}) \quad \sum_{i=0}^{N/2} (\pi_i)^{1/2} = O(N^{1/4})$$

419 and averaging over initial data $n_w(0) = j$, $j = 0, \dots, N/2$, we have

$$420 \quad (\text{A.10}) \quad E_\pi [\|\rho_m(j; i) - \pi_i\|_V] \leq \sum_{j=0}^{N/2} \pi_j \|\rho_m(j; \cdot) - \pi\|_V$$

$$421 \quad (\text{A.11}) \quad \leq \frac{1}{2} N^{-1/2} e^{-2c} \sum_{j=0}^{N/2} (\pi_j)^{1/2} \sum_{i=0}^{N/2} (\pi_i)^{1/2}$$

$$422 \quad (\text{A.12}) \quad \leq A e^{-2c}$$

423 for some A independent of N . This proves theorem 5.1.

424 **A.2. Lower bound.** In terms of the right-eigenvectors v_k , $k = 0, \dots, N/2$, (with
425 col sum =1)

$$426 \quad (\text{A.13}) \quad \Pr\{n_w(m) = i \mid n_w(0) = j\} = T_{ij}^m = \rho_m(i) \text{ if we take } j = 0$$

$$427 \quad (\text{A.14}) \quad = \sum_{k=0}^{N/2} \pi_i \lambda_k^m v_k(i) v_k(j)$$

$$428 \quad (\text{A.15}) \quad = \pi_i + \pi_i \sum_{k=0}^{N/2} \lambda_k^m v_k(i) v_k(j)$$

429 and

$$430 \quad (A.16) \quad E_{\rho_m}[v_1(i)] = \sum_{i=0}^{N/2} v_1(i) \rho_m(i) = \sum_{i=0}^{N/2} \pi_i v_1(i) \sum_{k=0}^{N/2} \lambda_k^m v_k(i) v_k(0)$$

$$431 \quad (A.17) \quad = \sum_{i=0}^{N/2} \sum_{k=0}^{N/2} \lambda_k^m \pi_i v_1(i) v_k(i) v_k(0)$$

$$432 \quad (A.18) \quad = \sum_{k=0}^{N/2} v_k(0) \lambda_k^m \sum_{i=0}^{N/2} \pi_i v_1(i) v_k(i)$$

$$433 \quad (A.19) \quad = v_1(0) \lambda_1^m \sum_{i=0}^{N/2} \pi_i v_1(i) v_1(i)$$

$$434 \quad (A.20) \quad = v_1(0) \lambda_1^m = v_1(0) \left(1 - \frac{4}{N}\right)^m$$

$$435 \quad (A.21) \quad E_{\pi}[v_1(i)] = 0; \text{var}_{\pi}\{v_1(i)\} = 1.$$

436 Next for $m = \frac{1}{8}N \log N - c\frac{N}{2}$, we get $E[v_1] = \frac{v_1(0)}{\sqrt{N}} e^{2c}$. A similar calculation gives

$$437 \quad (A.22) \quad E_{\rho_m}[v_2(i)] = \sum_{i=0}^{N/2} v_2(i) \rho_m(i) = \sum_{i=0}^{N/2} \pi_i v_2(i) \sum_{k=0}^{N/2} \lambda_k^m v_k(i) v_k(0)$$

$$438 \quad (A.23) \quad = \sum_{i=0}^{N/2} \pi_i v_2(i) \sum_{k=0}^{N/2} \lambda_k^m v_k(i) v_k(0) = \sum_{i=0}^{N/2} \sum_{k=0}^{N/2} \lambda_k^m \pi_i v_2(i) v_k(i) v_k(0)$$

$$439 \quad (A.24) \quad = \sum_{k=0}^{N/2} v_k(0) \lambda_k^m \sum_{i=0}^{N/2} \pi_i v_2(i) v_k(i) = v_2(0) \lambda_2^m \sum_{i=0}^{N/2} \pi_i v_2(i) v_2(i)$$

$$440 \quad (A.25) \quad = v_2(0) \lambda_2^m \sim v_2(0) \left(1 - \frac{8}{N}\right)^m,$$

$$441 \quad (A.26) \quad E_{\pi}[v_1(i)] = 0; \text{var}_{\pi}\{v_1(i)\} = 1$$

443 Next we deduce $v_1(0)$ and $v_2(0)$ from $v_1^2 = Av_2 + B$, $v_1(i) = C(N/4 - i)$, and the
444 orthogonality of v_i :

$$445 \quad (A.27) \quad 1 = \sum_{i=0}^{N/2} v_1(i)^2 \pi_i = \sum_{i=0}^{N/2} C^2(i - N/4)^2 \pi_i \sim C^2 \frac{N}{16}$$

446 Therefore, taking $C = \frac{4}{\sqrt{N}}$, we find $v_1(0) \sim \sqrt{N}$. Furthermore, $Av_2(0) = v_1(0) - b$.
447 Now, we have

$$448 \quad (A.28) \quad \text{Var}_{\rho_m}\{v_1\} = E_{\rho_m}[Av_2 + B] - N\lambda_1^{2m}$$

$$449 \quad (A.29) \quad = (N - B)\lambda_2^m + B - N\lambda_1^{2m} \sim B(1 - \lambda_2^m)$$

451 So with the same normalization as above for v_1 , we deduce $\text{Var}\{v_1'\}$ is uniformly
452 bounded by constant $2b$, since $B = b + O(\log N/N)$. Now, by Chebyshev's inequality,

$$453 \quad (A.30) \quad \text{Pr}_{\pi}\{|v_1| \leq k\} \geq 1 - \frac{1}{k^2}$$

454 and

$$\begin{aligned}
455 \quad (\text{A.31}) \quad & Pr_{\rho_m} \{|v_1| \leq k\} \leq Pr_{\rho_m} \{v_1 \leq k\} \\
456 \quad (\text{A.32}) \quad & = Pr_{\rho_m} \{E_{\rho_m}[v_1] - v_1 \geq E_{\rho_m}[v_1] - k\} \\
457 \quad (\text{A.33}) \quad & \leq Pr_{\rho_m} \{|E_{\rho_m}[v_1] - v_1| \geq |E_{\rho_m}[v_1] - k|\} \\
458 \quad (\text{A.34}) \quad & = Pr_{\rho_m} \{(E_{\rho_m}[v_1] - v_1)^2 \geq (E_{\rho_m}[v_1] - k)^2\} \\
459 \quad (\text{A.35}) \quad & \leq \frac{Var_{\rho_m}(v_1)}{(E_{\rho_m}[v_1] - k)^2} \\
460 \quad (\text{A.36}) \quad & \leq \frac{B}{(\sqrt{N}\lambda_1^m - k)^2} \\
461
\end{aligned}$$

462 Thus, if $K \subset \{0, \dots, N/2\}$ such that $|v_1| \leq k$ for $k \in K$, we deduce

$$\begin{aligned}
463 \quad (\text{A.37}) \quad & 2\|\rho_m - \pi\|_V = \sum_{i=0}^{N/2} |\rho_m(0, i) - \pi_i| \\
464 \quad (\text{A.38}) \quad & \geq \sum_K |\rho_m(0, i) - \pi_i| \\
465 \quad (\text{A.39}) \quad & \geq \sum_K \pi_i - \sum_K \rho_m(0, i) \\
466 \quad (\text{A.40}) \quad & \geq 1 - \frac{1}{k^2} - \frac{B}{(\sqrt{N}\lambda_1^m - k)^2} \\
467
\end{aligned}$$

468 Choose $k = d\sqrt{N}\lambda_1^m$ to obtain

$$469 \quad (\text{A.41}) \quad 2\|\rho_m - \pi\|_V \geq 1 - \left[\frac{1}{d^2} + \frac{B}{(1-d)^2} \right] N\lambda_1^{-2m}$$

$$470 \quad (\text{A.42}) \quad \geq 1 - be^{4c},$$

472 which proves theorem 5.2.

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